



LAZER DİYOT DENKLEMLERİNİN VOLTERRA SERİSİ İLE ÇÖZÜMLERİ

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VOLTERRA SERIES EXPANSION OF LASER DIODE RATE EQUATION

The basic laser diode equations are:

$$P' = \Gamma A(N - N_{tr})(1 - \hat{\epsilon}P)P - \frac{1}{\tau_p}P + \frac{\beta\Gamma N}{\tau_n} \quad (1)$$

$$N' = \frac{1}{qV}I - \frac{1}{\tau_n}N - \Gamma A[N - N_{tr}][1 - \hat{\epsilon}P]P \quad (2)$$

The expansion of equation (1) is:

$$\begin{aligned} P' &= \Gamma A[N - N_{tr} - \hat{\epsilon}PN + \hat{\epsilon}N_{tr}P]P - \frac{1}{\tau_p}P + \frac{\beta\Gamma N}{\tau_n} \\ &= \Gamma ANP - \Gamma AN_{tr}P - \Gamma A\hat{\epsilon}NP^2 + \Gamma A\hat{\epsilon}N_{tr}P^2 - \frac{1}{\tau_p}P + \frac{\beta\Gamma N}{\tau_n} \\ P' &= \frac{\beta\Gamma}{\tau_n}N - (\Gamma AN_{tr} + \frac{1}{\tau_p})P + \Gamma A\hat{\epsilon}N_{tr}P^2 + \Gamma ANP - \Gamma A\hat{\epsilon}NP^2 \end{aligned} \quad (3)$$

Adding equation (1) and (2)

$$P' + N' = \Gamma A(N - N_{tr})(1 - \hat{\epsilon}P)P - \frac{1}{\tau_p}P + \frac{\beta\Gamma N}{\tau_n} + \frac{1}{qV}I - \frac{1}{\tau_n}N - \Gamma A(N - N_{tr})(1 - \hat{\epsilon}P)P \quad (4)$$

$$P' + N' = \frac{I}{qV} - \frac{1}{\tau_n}(\beta\Gamma - 1)N - \frac{1}{\tau_p}P \quad (5)$$

The input current is composed of the sum of I_0 , D.C component, $\hat{I}$, a time-varying component. We thus let

$$I = I_0 + C\hat{I} \quad (6)$$

The Volterra expansion of the photon density and the electron density are:

$$P = \sum_{n=0}^{\infty} c^n P_n \quad \text{and} \quad N = \sum_{n=0}^{\infty} c^n N_n$$

substituting this expansion in equation (3) we obtain

$$\sum_{n=0}^{\infty} c^n P'_n = \frac{\beta\Gamma}{\tau_n} \sum_{n=0}^{\infty} c^n N_n - (\Gamma AN_{tr} + \frac{1}{\tau_p}) \sum_{n=0}^{\infty} c^n P_n +$$

$$\begin{aligned}
& + \Gamma A \hat{N}_{tr} \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} C^{n_1+n_2} P_{n_1} P_{n_2} \\
& + \Gamma A \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} C^{n_1+n_2} N_{n_1} P_{n_2} \\
& - \Gamma A \hat{\varepsilon} \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \sum_{n_3=0}^{\infty} C^{n_1+n_2+n_3} N_{n_1} P_{n_2} P_{n_3}
\end{aligned} \tag{7}$$

Yukarıdaki ifadenin acinini powerserisi turunden yapilir. Bu seri

$$\left[\sum_{n=0}^{\infty} C^n x_n \right]^2 = \sum_{n_1=0}^{\infty} C^{n_1} x_{n_1} \sum_{n_2=0}^{\infty} C^{n_2} x_{n_2} = \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} C^{n_1} x_{n_1} C^{n_2} x_{n_2} = \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} C^{n_1+n_2} x_{n_1} x_{n_2} \tag{8}$$

$$= a_0 c^0 + a_1 c^1 + a_2 c^2 + a_3 c^3 + a_4 c^4 + a_5 c^5 + \dots \tag{9}$$

$$n_1 + n_2 = 0 \rightarrow c^0$$

$$n_1 + n_2 = 1 \rightarrow c^1$$

$$n_1 + n_2 = 2 \rightarrow c^2$$

$$n_1 + n_2 = 3 \rightarrow c^3$$

$$n_1 + n_2 = 4 \rightarrow c^4$$

$$a_0 = x_0 x_0 = x^2$$

$$a_1 = x_0 x_1 + x_1 x_0 = 2x_0 x_1$$

$$a_1 = x_1 x_1 + x_0 x_2 + x_2 x_0 = x_1^2 + 2x_0 x_2$$

a_3	n_1	n_2	n_3
1	0	0	3
2	0	3	0
3	3	0	0
4	1	1	1
5	1	0	2
6	1	2	0
7	0	1	2
8	0	2	1
9	2	1	0
10	2	0	1
Degiskenler	N_{n1}	P_{n2}	P_{n3}

Terim sayisi ise

$$m = \frac{(n+k-1)!}{(k-1)n!} \tag{10}$$

$$m_{acinin(ise)} = \frac{(n+k-1)!}{(k-1)n!} = \frac{3+3-1}{(3-1)3!} = 10$$

$$\begin{aligned}
a_3 &= 2P_3 + 2P_0 + 2P_1 + 2P_2 \\
&N_0P_3 + P_0N_3 + N_1P_2 + N_2P_2 \\
&N_0P_0P_3 + N_0P_0P_3 + P_0P_0N_3 + N_1P_1P_1 + \\
&N_1P_0P_2 + N_1P_0P_2 + N_0P_1P_3 + N_0P_1P_2 + P_0P_1P_2 + P_0P_1P_2
\end{aligned}$$

$$\begin{aligned}
a_3 &= 2()P_3 + 2P_0 + 2P_1 + 2P_2 + \\
&N_0P_3 + P_0N_3 + N_1P_2 + N_2P_2 \\
&N_0P_0P_3 + N_0P_0P_3 + P_0P_0N_3 + N_1P_1P_1 + \\
&N_1P_0P_2 + N_1P_0P_2 + N_0P_1P_3 + N_0P_1P_2 + P_0P_1P_2 + P_0P_1P_2
\end{aligned}$$

we now equate like power of “c” to obtain

For c^0

$$P'_0 = \frac{\beta\Gamma}{\tau_n} N_0 - (\Gamma A N_{tr} + \frac{1}{\tau_p}) P_0 + \Gamma A \hat{N}_{tr} P_0^2 + \Gamma A N_0 P_0 - \Gamma A \hat{N}_0 P_0^2 \quad (4A)$$

For c^1

$$P'_1 = \frac{\beta\Gamma}{\tau_n} N_1 - (\beta\Gamma N_{tr} + \frac{1}{\tau_p}) P_1 + 2\Gamma A \hat{N}_{tr} P_0 P_1 + \Gamma A [N_0 P_1 + N_1 P_0 - \Gamma A \hat{N}_0 [2N_0 P_0 P_1 + P_0^2 N_1]] \quad (4B)$$

For c^2

$$\begin{aligned}
P'_2 &= \frac{\beta\Gamma}{\tau_n} N_2 - (\beta\Gamma N_{tr} + \frac{1}{\tau_p}) P_2 + \Gamma A \hat{N}_{tr} [2P_0 P_2 + P_1^2] + \Gamma A [N_0 P_2 + P_0 N_2 + N_1 P_1] \\
&- \Gamma A \hat{N}_0 [2N_0 P_0 P_2 + P_0^2 N_2 + N_0 P_1^2 + 2P_0 N_1 P_1]
\end{aligned} \quad (4C)$$

For c^3

$$\begin{aligned}
P'_3 &= \frac{\beta\Gamma}{\tau_n} N_3 - (\beta\Gamma N_{tr} + \frac{1}{\tau_p}) P_3 + \Gamma A \hat{N}_{tr} [2P_0 P_3 + 2P_1 P_2] + \Gamma A [N_0 P_3 + P_0 N_3 + N_1 P_2 + N_2 P_1] \\
&- \Gamma A \hat{N}_0 [2N_0 P_0 P_3 + P_0^2 N_3 + N_1 P_1^2 + 2N_1 P_0 P_2 + 2N_0 P_1 P_2 + 2P_0 P_1 N_2]
\end{aligned} \quad (4D)$$

Substituting the expansion for N and P in equation (3) we obtain

$$\sum_{n=0}^{\infty} c^n [P'_n + N'_n] = \frac{I_0}{qV} + c \frac{\hat{I}}{qV} + \frac{1}{\tau_n} (\beta\Gamma - 1) \sum_{n=0}^{\infty} c^n N_n - \frac{1}{\tau_p} \sum_{n=0}^{\infty} c^n P_n \quad (5)$$

For c^0

$$P'_0 + N'_0 = \frac{I_0}{qV} + \frac{1}{\tau_n}(\beta\Gamma - 1)N_0 - \frac{1}{\tau_p}P_0 \quad (5A)$$

For c^1

$$P'_1 + N'_1 = \frac{\hat{I}}{qV} + \frac{1}{\tau_n}(\beta\Gamma - 1)N_1 - \frac{1}{\tau_p}P_1 \quad (5B)$$

For c^2

$$P'_2 + N'_2 = \frac{1}{\tau_n}(\beta\Gamma - 1)N_2 - \frac{1}{\tau_p}P_2 \quad (5C)$$

For c^3

$$P'_3 + N'_3 = \frac{1}{\tau_n}(\beta\Gamma - 1)N_3 - \frac{1}{\tau_p}P_3 \quad (5D)$$

We now solve simultaneous equations (4) and (5) for the Volterra operator N_n and P_n successively.

To solve for N_0 and P_0 from equation (4A) and (5A) we note that $N'_0 + P'_0 = 0$ since $I_0 = \text{constant}$. Thus from equation(4A) we have

$$0 = \frac{\beta\Gamma}{\tau_n}N_0 - (\Gamma AN_r + \frac{1}{\tau_p})P_0 + \Gamma A\hat{e}N_rP_0^2 + \Gamma AN_0P_0 - \Gamma A\hat{e}N_0P_0^2 \quad (6A)$$

and from equation (5A)

$$0 = \frac{I_0}{qV} + \frac{1}{\tau_n}(\beta\Gamma - 1)N_0 - \frac{1}{\tau_p}P_0 \quad (6B)$$

solving equation(6B) for P_0 we have

$$P_0 = \frac{\tau_p}{qV}I_0 + \frac{\tau_p}{\tau_n}(\beta\Gamma - 1)N_0 \quad (6C)$$

now substitute equation (6C) into (6A) to obtain

$$\begin{aligned}
0 = & \frac{\beta\Gamma}{\tau_n} N_0 - (\Gamma A N_{tr} + \frac{1}{\tau_p}) \left[\frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \right] + \Gamma A \hat{\varepsilon} N_{tr} \left[\frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \right]^2 \\
& + \Gamma A N_0 \left[\frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \right] - \Gamma A \hat{\varepsilon} N_0 \left[\frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \right]^2
\end{aligned}$$

We now expand the brackets to obtain

$$\begin{aligned}
0 = & \frac{\beta\Gamma}{\tau_n} N_0 - (\Gamma A N_{tr} + \frac{1}{\tau_p}) \frac{\tau_p}{qV} I_0 \\
& - (\Gamma A \hat{\varepsilon} N_{tr} + \frac{1}{\tau_n}) \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \\
& + \Gamma A \hat{\varepsilon} N_{tr} \left[\frac{\tau_p}{qV} I_0 \right]^2 + \Gamma A \hat{\varepsilon} N_{tr} \left[\frac{\tau_p}{\tau_n} (\beta\Gamma - 1) \right]^2 N_0^2 \\
& + 2\Gamma A \hat{\varepsilon} N_{tr} \frac{\tau_p}{qV} I_0 \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \\
& + \Gamma A \frac{\tau_p I_0 N_0}{qV} + \left[\frac{\Gamma A \tau_p}{\tau_n} (\beta\Gamma - 1) \right] N_0^2 \\
& - \Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{qV} I_0 \right]^2 N_0 - \Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta\Gamma - 1) \right]^2 N_0^3 \\
& - 2\Gamma A \hat{\varepsilon} \frac{\tau_p}{qV} I_0 \frac{\tau_p}{\tau_n} (\beta\Gamma - 1)^2 N_0^2
\end{aligned}$$

Collecting terms in like power of N_0 we obtain the cubic:

$$\begin{aligned}
0 = & -\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta\Gamma - 1) \right]^2 N_0^3 \\
& + \left[\Gamma A \hat{\varepsilon} N_{tr} \left[\frac{\tau_p}{\tau_n} (\beta\Gamma - 1) \right]^2 + \frac{\Gamma A \tau_p}{\tau_n} (\beta\Gamma - 1) - 2\Gamma A \hat{\varepsilon} \frac{\tau_p}{qV} I_0 \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) \right] N_0^2 \\
& + \left[\frac{\beta\Gamma}{\tau_n} - (\Gamma A N_{tr} + \frac{1}{\tau_p}) \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) + 2\Gamma A \hat{\varepsilon} N_0 \frac{\tau_p}{qV} I_0 \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) + \Gamma A \frac{\tau_p}{qV} I_0 - \Gamma A \hat{\varepsilon} \left(\frac{\tau_p}{qV} I_0 \right)^2 \right] N_0 + \\
& \left[-(\Gamma A N_{tr} + \frac{1}{\tau_p}) \left(\frac{\tau_p}{qV} I_0 \right) + \Gamma A \hat{\varepsilon} N_{tr} \left(\frac{\tau_p}{qV} I_0 \right)^2 \right]
\end{aligned} \tag{6D}$$

This a cubic polynomial

$$N_0^3 + C_2 N_0^2 + C_1 N_0 + C_0 = 0 \quad (7)$$

in which the coefficients are:

$$C_2 = - \frac{\left[\Gamma A \hat{N}_{tr} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2 + \frac{\Gamma A \tau_p}{\tau_n} (\beta \Gamma - 1) - 2 \Gamma A \hat{\varepsilon} \frac{\tau_p^2}{qV} I_0 \frac{1}{\tau_n} (\beta \Gamma - 1) \right]}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2} \quad (7A)$$

$$C_1 = - \frac{\left[\frac{\beta \Gamma}{\tau_n} - \frac{\tau_p}{\tau_n} (\Gamma A N_{tr} + \frac{1}{\tau_p}) (\beta \Gamma - 1) + 2 \Gamma A \hat{N}_{tr} I_0 \frac{\tau_p^2}{qV \tau_n} (\beta \Gamma - 1) + \Gamma A \frac{\tau_p}{qV} I_0 - \Gamma A \hat{\varepsilon} \left(\frac{\tau_p}{qV} I_0 \right)^2 \right]}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2} \quad (7B)$$

$$C_0 = - \frac{\left[- (\Gamma A N_{tr} + \frac{1}{\tau_p}) \left(\frac{\tau_p}{qV} I_0 \right) + \Gamma A \hat{N}_{tr} \left(\frac{\tau_p}{qV} I_0 \right)^2 \right]}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2} \quad (7C)$$

Each of these coefficients is polynomial in I_0 .

Thus:

$$C_2 = c_{20} + c_{21} I_0 \quad (8)$$

in which

$$C_{20} = - \frac{\Gamma A \hat{N}_{tr} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2 + \Gamma A \frac{\tau_p}{\tau_n} (\beta \Gamma - 1)}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2}$$

$$C_{20} = - \frac{1 + \hat{\varepsilon} \frac{\tau_p}{\tau_n} (\beta \Gamma - 1) N_{tr}}{\hat{\varepsilon} \frac{\tau_p}{\tau_n} (\beta \Gamma - 1)} \quad (8A)$$

and

$$C_{21} = \frac{\frac{2\hat{\varepsilon}\tau_p^2}{qV\tau_n}(\beta\Gamma - 1)\Gamma A}{\Gamma A\hat{\varepsilon}\left[\frac{\tau_p}{\tau_n}(\beta\Gamma - 1)\right]^2} \quad (8B)$$

$$C_{21} = \frac{2\tau_n}{qV(\beta\Gamma - 1)} \quad (8B)$$

Also

$$C_1 = c_{10} + c_{11}I_0 + c_{12}I_0^2 \quad (9)$$

in which

$$C_{10} = -\frac{\frac{\beta\Gamma}{\tau_n} - \frac{\tau_p}{\tau_n}\left[\Gamma AN_{tr} + \frac{1}{\tau_p}\right](\beta\Gamma - 1)}{\Gamma A\hat{\varepsilon}\left[\frac{\tau_p}{\tau_n}(\beta\Gamma - 1)\right]^2}$$

$$C_{10} = \frac{1 - (\Gamma AN_{tr}\tau_p)(\beta\Gamma - 1)}{\Gamma A\hat{\varepsilon}\tau_p(\beta\Gamma - 1)^2} \frac{\tau_n}{\tau_p} \quad (9A)$$

$$C_{11} = -\frac{\frac{2\Gamma A\hat{\varepsilon}N_{tr}\tau_p^2}{qV\tau_n}(\beta\Gamma - 1)^2 + \frac{\Gamma A\tau_p}{qV\tau_n}}{\Gamma A\hat{\varepsilon}\left[\frac{\tau_p}{\tau_n}(\beta\Gamma - 1)\right]^2}$$

$$C_{11} = -\frac{2\hat{\varepsilon}N_{tr}\tau_p(\beta\Gamma - 1) + \tau_n}{qV\hat{\varepsilon}} \frac{\tau_n}{(\beta\Gamma - 1)^2 \tau_p} \quad (9B)$$

and

$$C_{12} = \frac{\Gamma A\hat{\varepsilon}\left[\frac{\tau_p}{qV}\right]^2}{\Gamma A\hat{\varepsilon}\left[\frac{\tau_p}{\tau_n}(\beta\Gamma - 1)\right]^2}$$

$$C_{12} = \left[\frac{\tau_n}{qV(\beta\Gamma - 1)}\right]^2 \quad (9C)$$

Also

$$C_0 = c_{01}I_0 + c_{02}I_0^2 \quad (10)$$

in which

$$C_{01} = \frac{\frac{\tau_p}{qV} \left(\Gamma A N_{tr} + \frac{1}{\tau_p} \right)}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} (\beta \Gamma - 1) \right]^2}$$

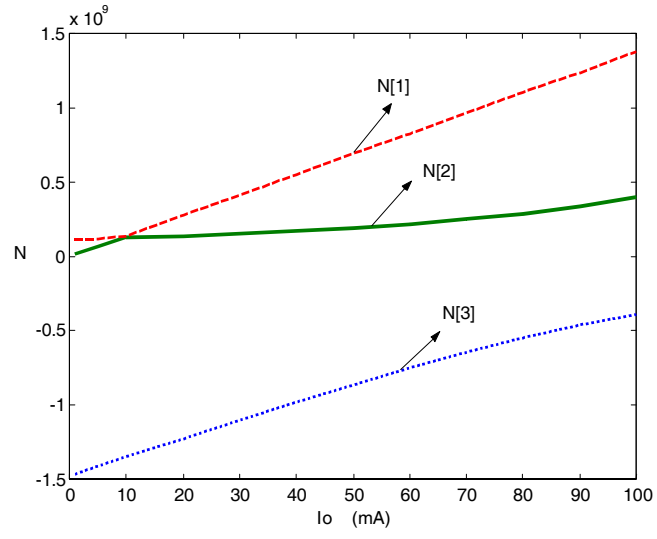
$$C_{01} = \frac{(1 + \Gamma A N_{tr} \tau_p) \tau_n^2}{q V \Gamma A \hat{\varepsilon} \tau_p^2 (\beta \Gamma - 1)^2} \quad (10A)$$

and

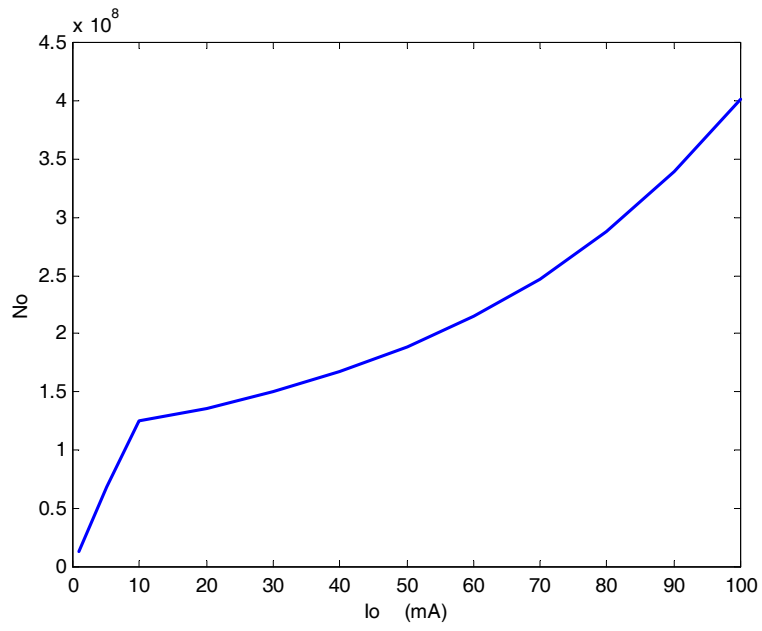
$$C_{02} = - \frac{\Gamma A \hat{\varepsilon} N_{tr} \left(\frac{\tau_p}{qV} \right)^2}{\Gamma A \hat{\varepsilon} \left[\frac{\tau_p}{\tau_n} \beta \Gamma - 1 \right]^2}$$

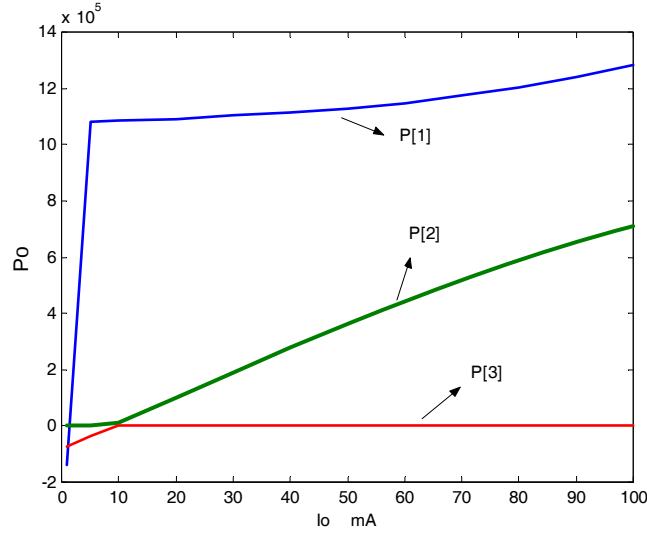
$$C_{02} = - \left[\frac{\tau_n}{qV(\beta \Gamma - 1)} \right]^2 N_{tr}$$

I[0] (mA)	N[1]	P[1]	N[2]	P[2]	N[3]	P[3]
1	-.1467259972e10	.140006e6	.13749985e8	.3102	.115969558e9	-.74338e5
5	-.1416062618e10	.1079832e7	.68747800e8	3.099	.119777690e9	-.37108e5
10	-.1352436817e10	.1083560e7	.124865398e9	.9191e4	.137538416e9	-24.938
15	-.1289270627e10	.1087623e7	.130477016e9	.55110e5	.206264734e9	-6.2158
20	-.1226622294e10	.1092062e7	.136583079e9	.100669e6	.275014463e9	-4.5184
30	-.1103158230e10	.1102272e7	.150624473e9	.190458e6	.412517255e9	-3.5491
40	-.982715730e9	.1114680e7	.167686571e9	.278049e6	.550020907e9	-3.2053
50	-.866183936e9	.1129933e7	.188659144e9	.362797e6	.687524790e9	-3.0291
60	-.754720709e9	.1148871e7	.214700197e9	.443859e6	.825028767e9	-2.9215
70	-.649776392e9	.1172550e7	.247260097e9	.520179e6	.962532794e9	-2.8495
80	-.553037478e9	.1202197e7	.288025379e9	.590533e6	.1100036849e10	-2.7985
90	-.466219398e9	.1239058e7	.338711478e9	.653671e6	.1237540920e10	-2.7591
100	-.390683388e9	.1284124e7	.400679635e9	.708605e6	.1375045003e10	-2.728
500	-.18366120e8	.5013356e7	.5528527641e10	.979373e6	.6875209739e10	-2.533
1000	-.2646556e7	.10001924e8	.1238801435e11	.990805e6	.1375041598e11.	-2.52
1238	-.3807e4	.12380000e8	.1565796988e11	.992727e6	.1702301405e11	-2.44
1239	.4970e4	.12389996e8	.1567171143e11	.992733e6	.1703676455e11	-2.51
1300	.512759e7	.12999627e8	.1650997879e11	.993102e6	.1787553974e11	-2.50
1500	.1867543e7	.14998641e8	.1925870654e11	.994088e6	.2062562220e11	-2.50



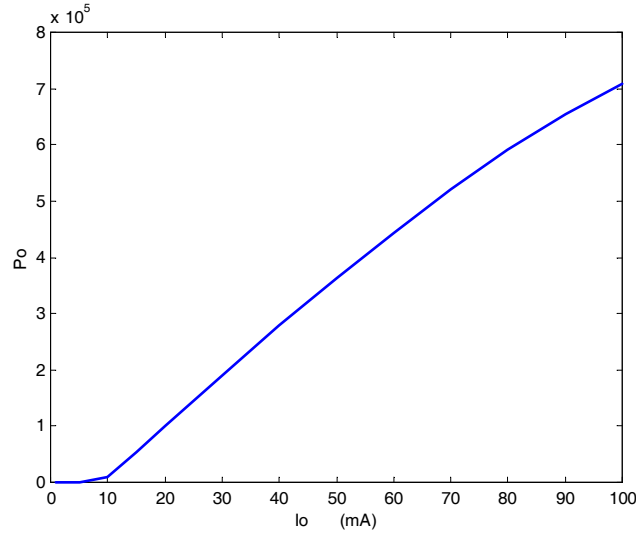
Sekil. Koklerin fiziksel durumları





There are three roots for N_0 of Equation (7). For each root, there is a corresponding value of P_0 obtained from equation (6C).

$$P_0 = \frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_0 \quad (6C)$$



for a physical solution, we require $N_0 > 0$ and $P_0 > 0$. We thus expect that only one of these solution pairs satisfy this condition for any value of I_0 . The procedure to determine N_0 and P_0 is to use equation (8), (9) and (10) to determine the coefficients C_0 , C_1 and C_2 as polynomials in I_0 and solving equation (7) and (6C) for the three solution pairs of N_0 and P_0 . If our equations are physically proper ones, there will be only one pair for which both N_0 and P_0 are positive and real. This is then the physical electron density N_0 , and photon density P_0 of the laser diode as a function of the injection current, I_0 .

Observe that for large values of I_0

$$C_2 \approx c_{21}I_0 = \frac{2\tau_n I_0}{qV(\beta\Gamma - 1)}$$

$$C_1 \approx c_{12}I_0^2 = \left(\frac{\tau_n}{qV(\beta\Gamma - 1)} \right)^2 I_0^2$$

$$C_0 \approx c_{02}I_0^2 = - \left(\frac{\tau_n}{qV(\beta\Gamma - 1)} \right)^2 N_{tr} I_0^2$$

For convenience, let

$$\alpha = \frac{\tau_n}{qV(\beta\Gamma - 1)}. \text{ Then}$$

$$C_2 \approx 2\alpha I_0$$

$$C_1 \approx \alpha^2 I_0^2$$

$$C_0 \approx -\alpha^2 I_0^2 N_{tr}$$

The cubic equation for N_0 , , equation(7), then is approximately

$$N_0^3 + 2\alpha N_0^2 I_0 + \alpha^2 I_0^2 N_0 - \alpha^2 I_0^2 N_{tr} = 0$$

For large values of I_0 one root of this equation is approximately obtained as

$$\alpha^2 I_0^2 N_0 - \alpha^2 I_0^2 N_{tr} = 0$$

For which the solution is

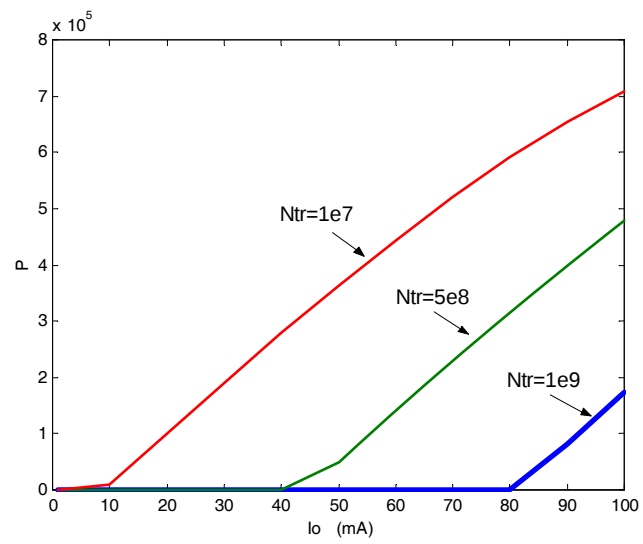
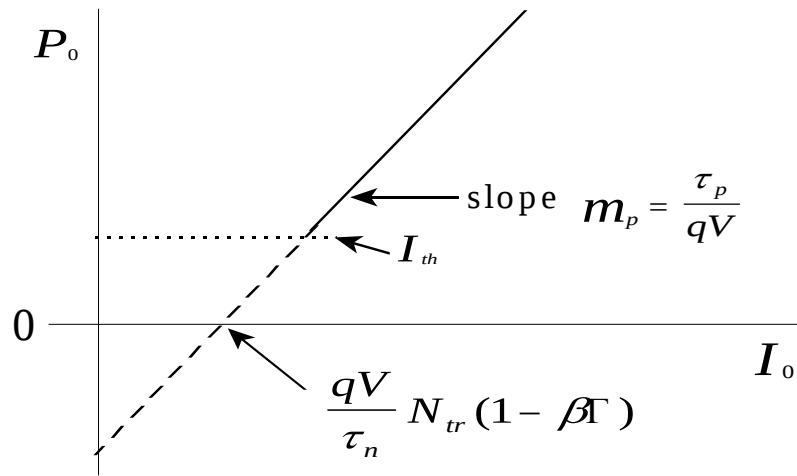
$$N_0 = N_{tr}$$

We thus observe that, for large values of I_0 , N_0 is a constant equal to N_{tr} . from a computer solution of equation (7), this is observed to be true for $I_0 \geq 7\text{ma}$.

We this approximation, the photon density for large ($I_0 \geq 7\text{ma}$) is from equation (6C)

$$P_0 \approx \frac{\tau_p}{qV} I_0 + \frac{\tau_p}{\tau_n} (\beta\Gamma - 1) N_{tr}$$

This the equation of a straight line as show



The slope of straight line is Figure

$$m_p = \frac{\tau_p}{qV}$$

Also, $P_0=0$ for

$$\frac{\tau_p}{qV} I_0 = \frac{\tau_p}{\tau_n} (1 - \beta\Gamma) N_{tr}$$

$$I_0 = \frac{qV}{\tau_n} (1 - \beta\Gamma) N_{tr} \dots \text{ampères}$$

this result is observed valid from an exact computer solution of the equation for $P_0 \geq 0$ note that we thus here that the threshold current $= I_{th}$

$$I_{th} = \frac{qV}{\tau_n} (1 - \beta\Gamma) N_{tr}. \quad (12)$$

This corresponds to the numerical value obtained experimentally. Note from our theoretical solution the threshold current, I_{th} is proportional to N_{tr} which is temperature dependent and increase with increasing temperature. However, the slope of $P_0 - I_0$ curves is $m_p = \tau_p / (qV)$ which is essentially independent of temperature. We thus expect the graph of $P_0 - I_0$ curves to be

We now determine the first-order Volterra operator for N_1 and P_1 by solving equation(4B) and (5B) simultaneously

$$P_1' = \frac{\beta\Gamma}{\tau_n} N_1 - \left(\Gamma A N_{tr} + \frac{1}{\tau_p} \right) P_1 + 2\Gamma A \hat{\varepsilon} N_{tr} P_0 P_1 + \Gamma A [N_0 P_1 + N_1 P_0] - \Gamma A \hat{\varepsilon} [2N_0 P_0 P_1 + P_0^2 N_1] \quad (4B)$$

$$P_1' + N_1' = \frac{1}{qV} \hat{I} + \frac{1}{\tau_n} (\beta\Gamma - 1) N_1 - \frac{1}{\tau_p} P_1 \quad (5C)$$

Collecting terms in equation(4B)

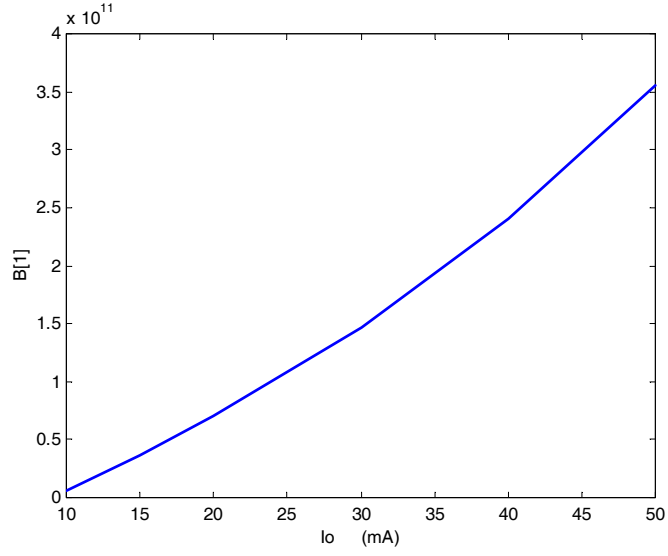
$$P_1' + \left[\Gamma A N_{tr} + \frac{1}{\tau_p} - 2\Gamma A \hat{\varepsilon} N_{tr} P_0 - \Gamma A N_0 + 2\Gamma A \hat{\varepsilon} N_0 P_0 \right] P_1 = \left[\frac{\beta\Gamma}{\tau_n} + \Gamma A P_0 - \Gamma A \hat{\varepsilon} P_0^2 \right] N_1 \quad (20)$$

This equation can be expressed as

$$P_1' + B_1 P_1 = B_0 N_1 \quad (21)$$

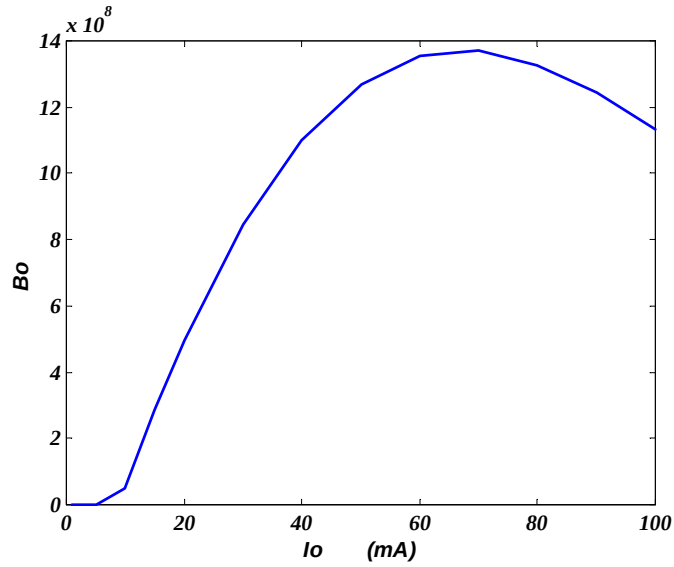
in which

$$B_1 = \Gamma A N_{tr} + \frac{1}{\tau_p} - 2\Gamma A \hat{N}_{tr} P_0 - \Gamma A N_0 + 2\Gamma A \hat{N}_0 P_0 \quad (21A)$$



and

$$B_0 = \frac{\beta \Gamma}{\tau_n} + \Gamma A P_0 - \Gamma A \hat{P}_0^2 \quad (21B)$$



Equation (21) can be solved for N₁ as

$$N_1 = \frac{1}{B_0} P_1' + \frac{B_1}{B_0} P_1 \quad (21C)$$

We now substitute equation (21C) into equation (5B)

$$P_1' + \frac{1}{B_0} P_1'' + \frac{B_1}{B_0} P_1' = \frac{1}{qV} \hat{I} + \frac{1}{\tau_n} \left(\frac{\beta\Gamma - 1}{B_0} \right) [P_1' + B_1 P_1] - \frac{1}{\tau_p} P_1$$

collecting terms

$$\frac{1}{B_0} P_1'' + \left(1 + \frac{B_1}{B_0} - \frac{1}{\tau_n} \frac{\beta\Gamma - 1}{B_0} \right) P_1' + \left(\frac{1}{\tau_p} - \frac{1}{\tau_n} \frac{\beta\Gamma - 1}{B_0} \right) P_1 = \frac{1}{qV} \hat{I}$$

multiplying by B_0

$$P_1'' + \left(B_0 + B_1 - \frac{\beta\Gamma - 1}{\tau_n} \right) P_1' + \left(\frac{B_0}{\tau_p} - \frac{(\beta\Gamma - 1)B_1}{\tau_n} \right) P_1 = \frac{B_0}{qV} \hat{I} \quad (22)$$

Equation (22) is a linear differential equation

$$P_1'' + D_1 P_1' + D_0 P_1 = \frac{B_0}{qV} \hat{I} \quad (23)$$

In which

$$D_1 = B_0 + B_1 - \left(\frac{\beta\Gamma - 1}{\tau_n} \right) \quad (23A)$$

and

$$D_0 = \frac{B_0}{\tau_p} - \left(\frac{\beta\Gamma - 1}{\tau_n} \right) B_1 \quad (23B)$$

The Laplace transform of this equation is

$$[s^2 + D_1 s + D_0] P_1(s) = \frac{B_0}{qV} \hat{I}(s)$$

or

$$P_1(s) = \frac{B_0}{qV} \left(\frac{1}{s^2 + D_1 s + D_0} \right) \hat{I}(s)$$

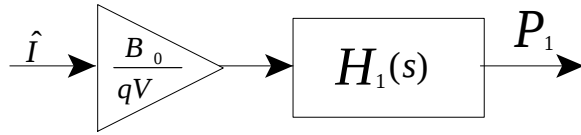
Then

$$P_1(s) = \frac{B_0}{qV} H_1(s) \hat{I}(s)$$

in which

$$H_1(s) = \frac{1}{s^2 + D_1 s + D_0}$$

which can be represented as

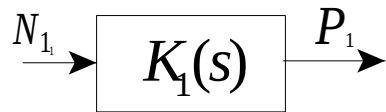


Figure

Now the Laplace transform of equation (21C) is

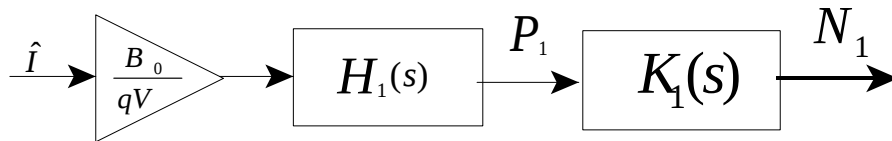
$$N_1(s) = \frac{1}{B_0} [s + B_1] P_1(s)$$

$$N_1(s) = K_1(s) P_1(s)$$



Figure

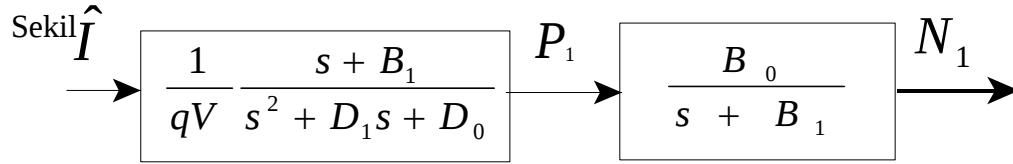
We thus have



Figure

$$H_1(s)K_1(s) = \frac{1}{qV} \frac{s + B_1}{s + D_1s + D_0}$$

So also



Figure

We now examine the constants D_0 and D_1 substituting equation (21B) into equation (23A), we obtain

$$D_1 = \Gamma A N_{tr} + \frac{1}{\tau_p} - 2\Gamma A \hat{N}_{tr} P_0 - \Gamma A N_0 + 2\Gamma A \hat{N}_0 P_0 + \frac{\beta\Gamma}{\tau_n} + \Gamma A P_0 - \Gamma A \hat{P}_0^2 - \frac{\beta\Gamma - 1}{\tau_n}$$

Collecting terms we have obtain

$$D_1 = \Gamma A (N_{tr} - N_0) + 2\Gamma A \hat{P}_0 (N_0 - N_{tr}) + \Gamma A (P_0 - \hat{P}_0^2) + \frac{1}{\tau_n} + \frac{1}{\tau_p}$$

$$D_1 = \Gamma A (1 - 2\hat{P}_0) (N_{tr} - N_0) + \Gamma A P_0 (1 - \hat{P}_0^2) + \frac{1}{\tau_n} + \frac{1}{\tau_p} \quad (24A)$$

Similarly for D_0 we substitute equation (21A) and (21B) into equation (23B) to obtain

$$\begin{aligned} D_0 = & \frac{\beta\Gamma}{\tau_n \tau_p} + \frac{\Gamma A}{\tau_p} P_0 - \frac{\Gamma A \hat{P}_0^2}{\tau_p} - \frac{\beta\Gamma - 1}{\tau_n} \Gamma A N_{tr} \\ & - \frac{\beta\Gamma - 1}{\tau_n \tau_p} + \frac{2\Gamma A \hat{N}_{tr} (\beta\Gamma - 1)}{\tau_n} P_0 \\ & + \frac{\beta\Gamma - 1}{\tau_n} \Gamma A N_0 - \frac{2\Gamma A \hat{N}_0 P_0 (\beta\Gamma - 1)}{\tau_n} \end{aligned} \quad (24A)$$

collecting terms,

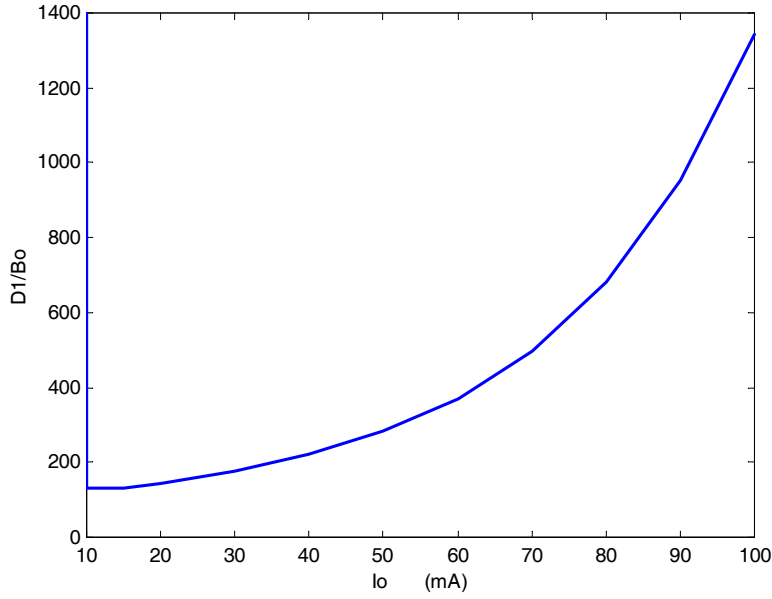
$$D_0 = \frac{\beta\Gamma - 1}{\tau_n} \Gamma A [N_0 - N_{tr}] + \frac{2\Gamma A \hat{\varepsilon} (\beta\Gamma - 1) P_0}{\tau_n} [N_{tr} - N_0] + \frac{\Gamma A}{\tau_p} P_0 - \frac{\Gamma A \hat{\varepsilon}}{\tau_p} P_0^2 + \frac{1}{\tau_n \tau_p}$$

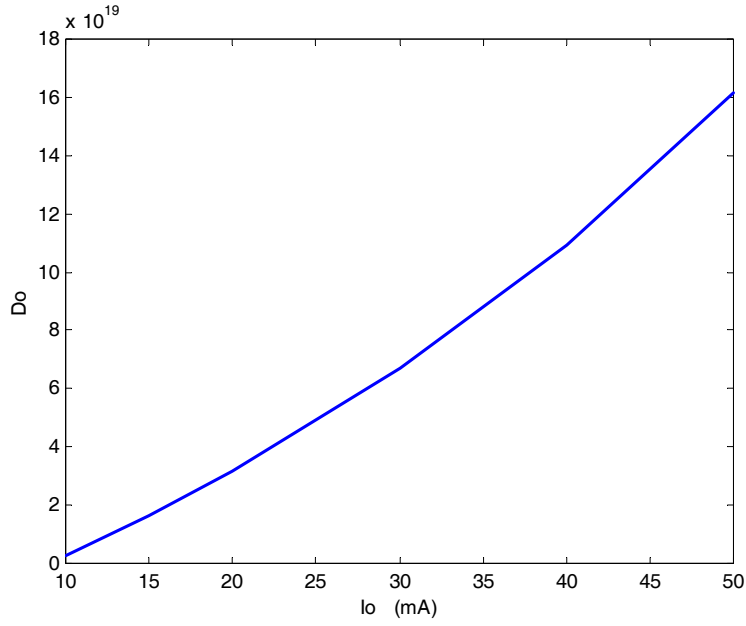
$$D_0 = \frac{\beta\Gamma - 1}{\tau_n} \Gamma A (1 - 2\hat{\varepsilon} P_0) [N_0 - N_{tr}] + \frac{\Gamma A}{\tau_p} P_0 (1 - \hat{\varepsilon} P_0^2) + \frac{1}{\tau_n \tau_p} \quad (24B)$$

We now observe that for large I_0 , $N_0 = N_{tr}$. Thus for large I_0

$$D_1 = \Gamma A P_0 (1 - \hat{\varepsilon} P_0) + \frac{1}{\tau_n} + \frac{1}{\tau_p} \quad (25A)$$

$$D_0 = \frac{\Gamma A}{\tau_p} P_0 (1 - \hat{\varepsilon} P_0) + \frac{1}{\tau_n} \frac{1}{\tau_p} \quad (25B)$$





The poles of $H_1(s)$ are the roots of

$$s^2 + D_1 s + D_0 = \left(s + \frac{D_1}{2}\right)^2 - \left(\frac{D_1^2}{4} - D_0\right)$$

$$= s + \frac{D_1}{2} \pm \sqrt{\frac{D_1^2}{4} - D_0}$$

so the roots are

$$s_1 = -\frac{D_1}{2} + \sqrt{\left(\frac{D_1}{2}\right)^2 - D_0}$$

$$s_2 = -\frac{D_1}{2} - \sqrt{\left(\frac{D_1}{2}\right)^2 - D_0}$$

We thus have

$$s_{1,2} = -\frac{D_1}{2} + \left[1 \pm \sqrt{1 - \frac{4D_0}{D_1^2}}\right]$$

now

$$\frac{4D_0}{D_1^2} < 10^{-18} \text{ so that}$$

$$\sqrt{1 - \frac{4D_0}{D_1^2}} \approx 1 - \frac{2D_0}{D_1^2}$$

Consequently

$$s_{1,2} = -\frac{D_1}{2} \left[1 \pm \left(1 - \frac{2D_0}{D_1^2} \right) \right]$$

or

$$s_1 \approx -D_1$$

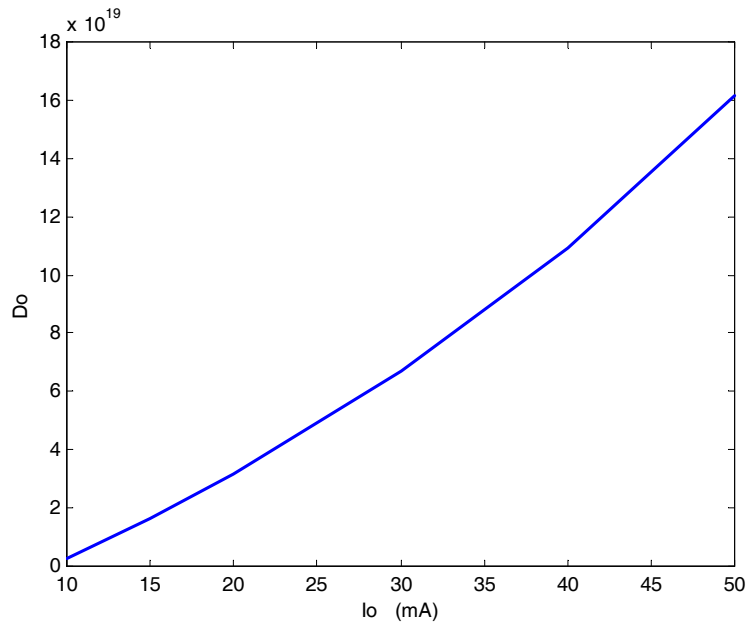
$$s_2 = -\frac{D_1}{2} \left(\frac{2D_0}{D_1^2} \right) = -\frac{D_0}{D_1}$$

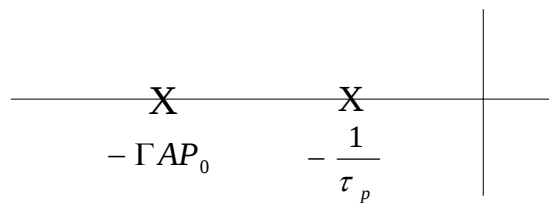
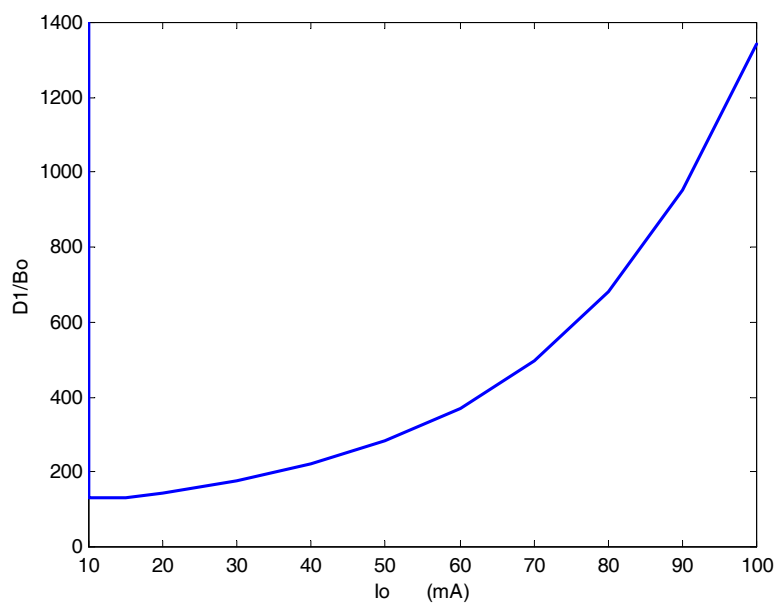
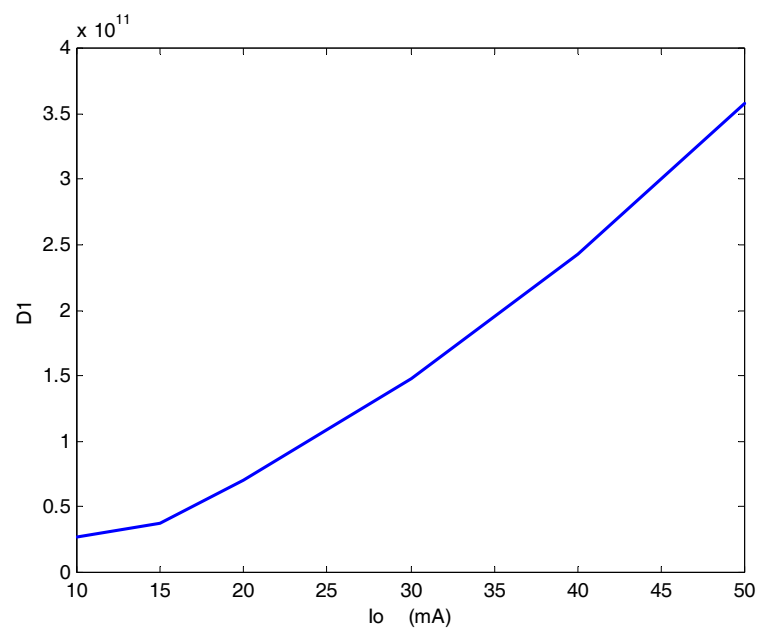
note that if $1 - 2\hat{\mathcal{E}}P_0 \approx 1$, we would have

$D_1 \approx \Gamma A P_0$ and $D_1 \approx (\Gamma A P_0)/\tau_p$ so that

$$s_1 \approx -\Gamma A P_0$$

$$s_2 \approx -\frac{1}{\tau_p}$$





Figure

Note that the system is stable. Also, are pole is proportional the P_0 which is proportional to I_0 .

We now determine the second-order Volterra operator, P_2 and N_2 . For this We use equation(4C) and (5C)

$$P_2' = \frac{\beta\Gamma}{\tau_n} N_2 - \left(\beta\Gamma N_{tr} + \frac{1}{\tau_p}\right) P_2 + \Gamma A \hat{N}_{tr} [2P_0 P_2 + P_1^2] + \Gamma A [N_0 P_2 + P_0 N_2 + N_1 P_1] - \Gamma A \hat{P}_0 [2N_0 P_0 P_2 + P_0^2 N_2 + N_0 P_1^2 + 2P_0 N_1 P_1] \quad (4C)$$

$$P_2' + N_2' = \frac{1}{\tau_n} (\beta\Gamma - 1) N_2 - \frac{1}{\tau_p} P_2 \quad (5C)$$

Collecting terms in equation(4C)

$$P_2' + \left[\Gamma A N_{tr} + \frac{1}{\tau_p} - 2\Gamma A \hat{N}_{tr} P_0 - \Gamma A \hat{N}_0 + 2\Gamma A \hat{N}_0 P_0 \right] P_2 = \left(\frac{\beta\Gamma}{\tau_n} + \Gamma A P_0 - \Gamma A \hat{P}_0^2 \right) N_2 + \left(\Gamma A \hat{N}_{tr} P_1^2 + \Gamma A N_1 P_1 - \Gamma A \hat{N}_0 P_1^2 - 2\Gamma A \hat{P}_0 N_1 P_1 \right) \quad (40)$$

This equation can be expressed as

$$P_2' + B_1 P_2 = B_0 N_2 + F_2 \quad (41)$$

in which B_0 and B_1 are given by equation (21B) and (21A) respectively and

$$F_2 = \Gamma A \hat{N}_{tr} P_1^2 + \Gamma A N_1 P_1 - \Gamma A \hat{N}_0 P_1^2 - 2\Gamma A \hat{P}_0 N_1 P_1$$

$$F_2 = \Gamma A N_1 P_1 [1 - 2\hat{P}_0] - \Gamma A \hat{P}_1^2 [N_0 - N_{tr}] \quad (41A)$$

Equation 41 can be solved for N_2 as

$$N_2 = \frac{1}{B_0} P_2' + \frac{B_1}{B_0} P_2 - \frac{1}{B_0} F_2 \quad (41B)$$

We now substitute equation (41B) in equation (5C)

$$P_2' + \frac{1}{B_0} P_2'' + \frac{B_1}{B_0} P_2' - \frac{1}{B_0} F_2' = \frac{(\beta\Gamma - 1)}{\tau_n} \frac{1}{B_0} P_2' + \frac{(\beta\Gamma - 1)}{\tau_n} \frac{B_1}{B_0} P_2 - \frac{(\beta\Gamma - 1)}{\tau_n B_0} F_2 - \frac{1}{\tau_p} P_2$$

Collecting terms

$$\frac{1}{B_0} P_2'' + \left[1 + \frac{B_1}{B_0} - \frac{(\beta\Gamma - 1)}{\tau_n} \frac{1}{B_0} \right] P_2' + \left[\frac{1}{\tau_p} - \frac{(\beta\Gamma - 1)}{\tau_n} \frac{B_1}{B_0} \right] P_2 = \frac{1}{B_0} F_2' - \frac{(\beta\Gamma - 1)}{\tau_n B_0} F_2$$

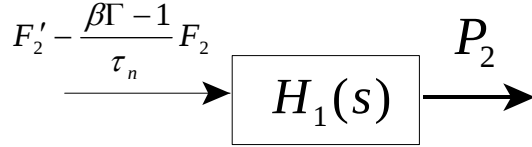
Multiplying by B_0

$$P_2'' + \left[B_0 + B_1 - \frac{(\beta\Gamma - 1)}{\tau_n} \right] P_2' + \left[\frac{B_0}{\tau_p} - \frac{(\beta\Gamma - 1)}{\tau_n} B_1 \right] P_2 = F_2' - \frac{(\beta\Gamma - 1)}{\tau_n} F_2 \quad (42)$$

This equation is a linear differential equation equation

$$P_2'' + D_1 P_2' + D_0 P_2 = F_2' - \frac{(\beta\Gamma - 1)}{\tau_n} F_2 \quad (43)$$

In which D_1 and D_0 are given by equation (23A) and (23B) respectively. A block diagram thus in



Figure

Now from equation (41B)

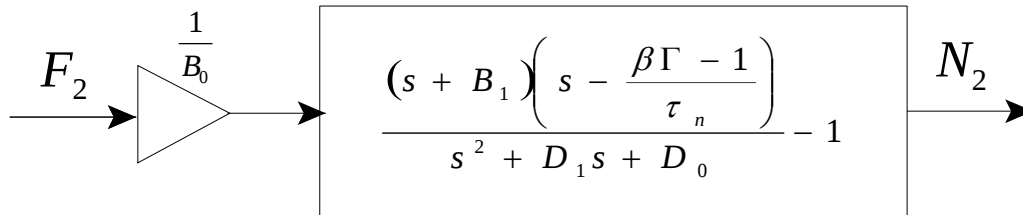
$$N_2(s) = \frac{s + B_1}{B_0} P_2(s) - \frac{1}{B_0} F_2(s)$$

and from equation (43)

$$P_2(s) = \frac{s - \frac{\beta\Gamma - 1}{\tau_n}}{s^2 + D_1 s + D_0} F_2(s)$$

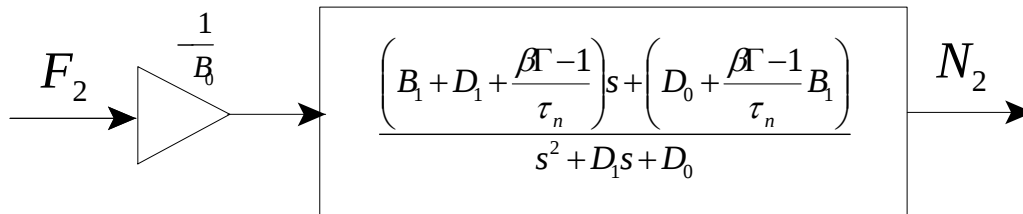
so

$$N_2(s) = \frac{1}{B_0} \left[\frac{(s + B_1) \left(s - \frac{\beta \Gamma - 1}{\tau_n} \right)}{s^2 + D_1 s + D_0} - 1 \right] F_2(s)$$



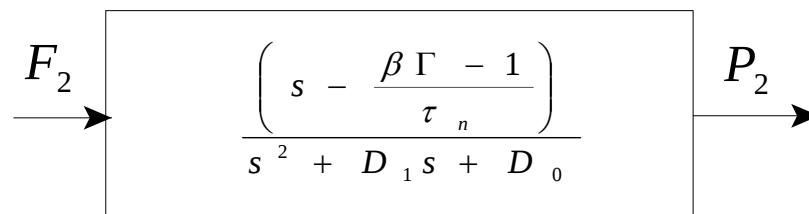
Figure

Which is equivalent to



Figure

and



Figure

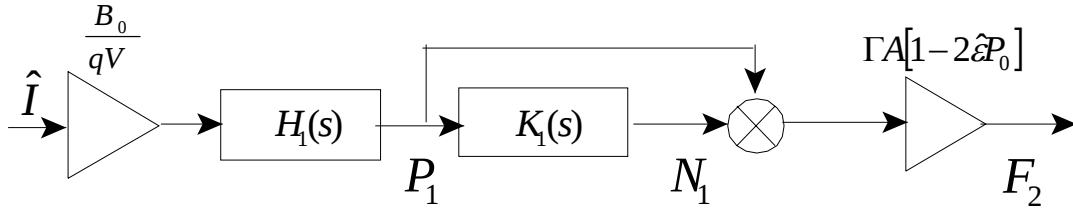
Now from Equation (41A)

$$F_2 = \Gamma A N_1 P_1 [1 - 2\hat{\mathcal{P}}_0] - \Gamma A \hat{\mathcal{P}}_1^2 [N_0 - N_{tr}]$$

For large I_0 , $N_0 = N_{tr}$ so that

$$F_2 = \Gamma A N_1 P_1 [1 - 2\hat{\mathcal{P}}_0]$$

which is obtained as



Figure

We now determine the third-order Volterra operators, P_3 and N_3 . For this, We use equation (4D) and (5D)

$$P'_3 = \frac{\beta\Gamma}{\tau_n} N_3 - \left(A\Gamma N_{tr} + \frac{1}{\tau_p} \right) P_3 + \Gamma A \hat{\mathcal{N}}_{tr} [2P_0 P_3 + 2P_1 P_2] + \Gamma A [N_0 P_3 + P_0 N_3 + N_1 P_2 + N_2 P_1] - \Gamma A \hat{\mathcal{E}} [2N_0 P_0 P_3 + P_0^2 N_3 + N_1 P_1^2 + 2N_1 P_0 P_2 + 2N_0 P_1 P_2 + 2P_0 P_1 N_2] \quad (4D)$$

$$P'_3 + N'_3 = \frac{1}{\tau_n} (\beta\Gamma - 1) N_3 - \frac{1}{\tau_p} P_3 \quad (5D)$$

Collecting terms in equation (4D)

$$P'_3 + \left[\Gamma A N_{tr} + \frac{1}{\tau_p} - 2\Gamma A \hat{\mathcal{N}}_{tr} P_0 - \Gamma A N_0 + 2\Gamma A \hat{\mathcal{N}}_0 P_0 \right] P_3 = \left[\frac{\beta\Gamma}{\tau_p} + \Gamma A P_0 - \Gamma A \hat{\mathcal{P}}_0^2 \right] N_3 + \left[2\Gamma A \hat{\mathcal{N}}_{tr} P_1 P_2 + \Gamma A N_1 P_2 + \Gamma A N_2 P_1 - \Gamma A \hat{\mathcal{N}}_1 P_1^2 - 2\Gamma A \hat{\mathcal{N}}_1 P_0 P_2 - 2\Gamma A \hat{\mathcal{N}}_0 P_1 P_2 - 2\Gamma A \hat{\mathcal{N}}_2 P_0 P_1 \right] \quad (60)$$

This equation can be expressed as

$$P'_3 + B_1 P_3 = B_0 N_3 + F_3 \quad (61)$$

in which B_0 and B_1 are given by equation (21B) and (21A) respectively and

$$\begin{aligned} F_3 &= \left[2\Gamma A \hat{N}_r P_1 P_2 + \Gamma A N_1 P_2 + \Gamma A N_2 P_1 - \Gamma A \hat{N}_1 P_1^2 - 2\Gamma A \hat{N}_1 P_0 P_2 - 2\Gamma A \hat{N}_0 P_1 P_2 - 2\Gamma A \hat{N}_2 P_0 P_1 \right] \\ &= 2\Gamma A \hat{P}_1 P_2 [N_r - N_0] + \Gamma A N_1 P_2 [1 - 2\hat{P}_0] + \Gamma A N_2 P_1 [1 - 2\hat{P}_0] - \Gamma A \hat{N}_1 P_1^2 \end{aligned}$$

Equation (61) can be solved for N_3 as

$$N_3 = \frac{1}{B_0} P_3' + \frac{B_1}{B_0} P_3 - \frac{1}{B_0} F_3 \quad (61B)$$

We now substitute equation (61B) in equation (5D)

$$P_3' + \frac{1}{B_0} P_3'' + \frac{B_1}{B_0} P_3' + \frac{1}{B_0} F_3' = \frac{(\beta\Gamma - 1)}{\tau_n} \frac{1}{B_0} P_3' + \frac{(\beta\Gamma - 1)}{\tau_n} \frac{B_1}{B_0} P_3' - \frac{(\beta\Gamma - 1)}{\tau_n} \frac{1}{B_0} F_3' - \frac{1}{\tau_p} P_3$$

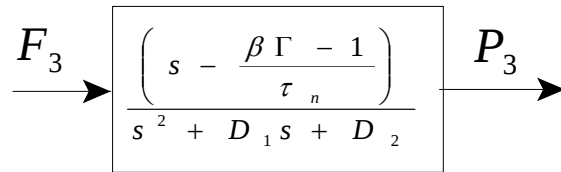
Collecting terms and multiplying by B_0 we obtain

$$P_3'' + \left[B_0 + B_1 - \frac{(\beta\Gamma - 1)}{\tau_n} \right] P_3' + \left[\frac{B_0}{\tau_p} - \frac{(\beta\Gamma - 1)}{\tau_n} B_1 \right] P_3 = F_3' + \frac{(\beta\Gamma - 1)}{\tau_n} F_3 \quad (62)$$

This is a linear differential equation

$$P_3'' + D_1 P_3' + D_0 P_3 = F_3' + \frac{(\beta\Gamma - 1)}{\tau_n} F_3 \quad (63)$$

In which D_1 and D_0 are given by equation (23A) and (23B) respectively. Similar to second-order operator, we have



Figure

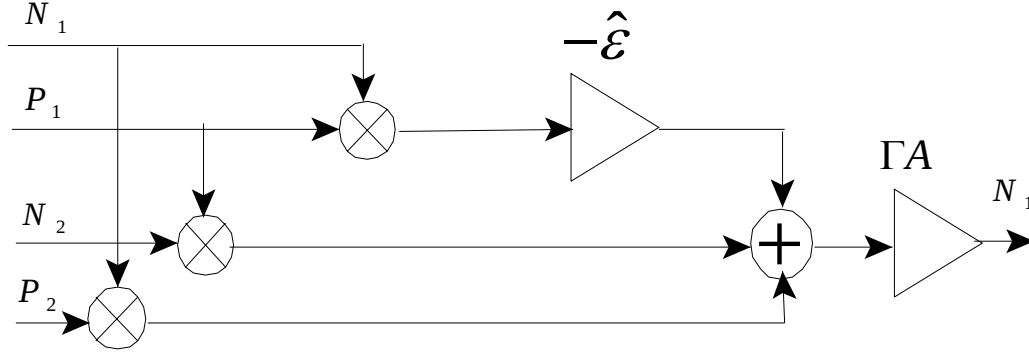
In which F_3 is given by equation (61A) note that $N_0 = N_{tr}$ for large I_0 so that for large I_0

$$F_3 = \Gamma A N_1 P_2 [1 - 2\hat{P}_0] + \Gamma A N_2 P_1 [1 - 2\hat{P}_0] - \Gamma A \hat{N}_1 P_1^2$$

Also, for $1 - 2\hat{\varepsilon}P_0 \approx 1$ we have

$$F_3 = \Gamma A N_1 P_2 + \Gamma A N_2 P_1 - \Gamma A \hat{\varepsilon} N_1 P_1^2 = \Gamma A (N_1 P_2 + N_2 P_1 - \hat{\varepsilon} N_1 P_1^2)$$

which is obtained as



Figure

In which N_1 , P_1 , N_2 , and P_2 are first and second-order operator of n and P respectively. Also N_3 , if desired, is obtained from equation (61B).

For linear feedback in which $K=K$,

$$P_1 = H_1 R_1 \quad (116)$$

$$P_2 = R_1 H_2 R_1 \quad (117)$$

$$P_3 = R_1 H_3 R_1 - 2R_1 H_2 \{I, R_1, K_1, H_2\} R_1$$

$$P_3 = R_1 H_3 R_1 + R_1 H_2 [I - R_1 K_1 H_2] R_1 - R_1 H_2 R_1 K_1 H_2 R_1 - R_1 H_2 R_1 \quad (118)$$

$$P_3 = R_1 [H_3 + H_2 [I - R_1 K_1 H_2] - H_2 R_1 K_1 H_2 - H_2] R_1$$

see too #7 for linearization using NL feedback in our case,

$$K_1 = K e^{-st_0} = K e^{-s_n} \quad (119)$$

note that feedback system operator, P_n are stable operators if R_1 is a stable operator now

$$R_1(s) = \frac{1}{1 + H_1(s) K_1(s)} \quad (120)$$

$$R_1(s) = \frac{1}{1 + \frac{B_0 K}{q} \frac{e^{-st_0}}{s^2 + D_1 s + D_0}} \quad (121)$$

$$R_1(s) = \frac{s^2 + D_1 s + D_0}{s^2 + D_1 s + D_0 + \frac{B_0 K}{q} e^{-st_0}} \quad (122)$$

The eq (27) is a stable operator if the roots of

$$F(s) = s^2 + D_1 s + D_0 + \frac{B_0 K}{q} e^{-t_0 s} \quad (123)$$

are all in the left- half plane.

For convenience, we define $s_n = t_0 s = \sigma_n + j\omega_n$ in which $\sigma_n = t_0 \sigma$ and $\omega_n = t_0 \omega$. Then

$$F(s_n) = s_n^2 + D_1 t_0 s_n + D_0 t_0^2 + \frac{B_0 K t_0^2}{q} e^{-s_n} \cos \omega_n \quad (124)$$

Then

$$F_r(\sigma_n, \omega_n) = \sigma_n^2 - \omega_n^2 + D_1 t_0 \sigma_n + D_0 t_0^2 + \frac{B_0 K t_0^2}{q} e^{-s_n} \cos \omega_n \quad (125)$$

$$F_i(\sigma_n, \omega_n) = 2\sigma_n - \omega_n + D_1 t_0 \sigma_n - \frac{B_0 K t_0^2}{q} e^{-s_n} \sin \omega_n \quad (126)$$

if $F(s_n) = 0$ then $F(s_n) = F_r(\sigma_n, \omega_n) + jF_i(\sigma_n, \omega_n) = 0$. Now $F=0$ if only if F_r and $F_i=0$. From eq.(125)

$$F_i(\sigma_n, \omega_n) = 0$$

$$\frac{\sin \omega_n}{\omega_n} = \frac{q}{B_0 K t_0^2} [D_1 t_0 + 2\sigma_n] e^{+\sigma_n} \quad (127)$$

To the equation to be stable, this equation can not have any solutions for $\sigma_n \geq 0$. Now the right side of eq. (127) is a monotonically increasing function of σ_n . Thus, eq(127) has no solutions for $\sigma_n \geq 0$ if

$$\frac{q D_1 t_0}{B_0 K t_0^2} = \frac{q D_1}{B_0 K t_0} > 1 \quad (128)$$

or

$$K < \frac{q D_1}{B_0 t_0} = K_0 \quad (129)$$

The range of K which the system is stable can be large since we have not considered that eq(125) must be simultaneously satisfied. For our laser diode parameters, result of a computer simulate are

t_0	10^{-12}	10^{-11}	10^{-10}
$K_{\max}/K_0$	1.01	1.26	1.32

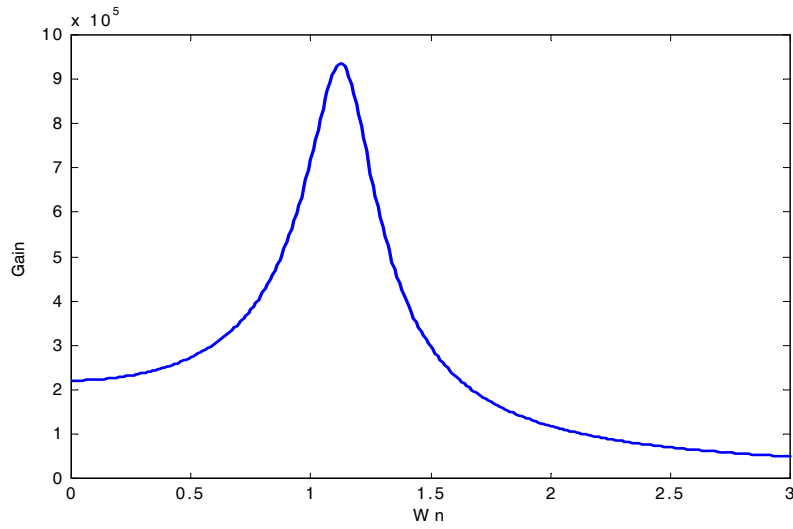
In which $K_{\max}$ is the maximum value of K for which the equation is stable. We thus observe that K_0 is a reasonable lower bound of K for stability.

For $K_0 \geq K \geq 0$, the system definitely is stable. For $K \in [0, K_0]$ we than have

$$P_1(j\omega_n) = \frac{B_0}{q} \frac{1}{(D_0 t_0^2 - \omega_n^2) + jD_1 t_0 \omega_n + \frac{B_0 K}{q} e^{-j\omega_n}} \quad (130)$$

and

$$R_1(j\omega_n) = \frac{(D_0 t_0^2 - \omega_n^2) + jD_1 t_0 \omega_n}{(D_0 t_0^2 - \omega_n^2) + jD_1 t_0 \omega_n + \frac{B_0 K}{q} e^{-j\omega_n}} \quad (131)$$



VOLTERRA ANALİZİ EŞİTLİKLERİ

$$\dot{I}_1 = A_1 \sin (wt + \phi_1)$$

$$\begin{aligned} \dot{I}_2 = \dot{I}_1^2 &= A_1^2 \sin^2 (wt + \phi_1) \\ &= \frac{A_1^2}{2} - \frac{A_1^2}{2} \cos(2wt + 2\phi_1) \end{aligned}$$

$$\dot{I}_3 = C_2 \dot{I}_2 = \frac{A_1^2}{2} - \frac{A_1^2}{2} \cos[2wt + 2\phi_1]$$

$$\dot{I}_4 = A_1 |F_3| \sin(wt + \phi_1 + \phi_3)$$

$$|F_3| = G_{33} \quad \phi_3 = \phi_{31}$$

$$\dot{I}_4 = A_1 G_{33} \sin(wt + \phi_1 + \phi_{31})$$

$$\dot{I}_4 = A_2 \sin(wt + \phi_1 + \phi_{31})$$

$$\begin{aligned} \phi_2 &= \phi_1 + \phi_{31} \\ A_2 &= A_1 G_{33} \end{aligned}$$

$$\dot{I}_4 = A_1 G_{33} \sin(wt + \phi_1)$$

$$\dot{I}_5 = \dot{I}_4 \dot{I}_1$$

$$= A_1^2 G_{33} \sin(wt + \phi_2) \sin(wt + \phi_1)$$

$$= \frac{A_1^2 G_{33}}{2} \cos(\phi_2 - \phi_1) - \frac{A_1^2 G_{33}}{2} \cos[2wt + \phi_2 + \phi_1]$$

$$= \frac{A_1^2 G_{33}}{2} \cos(\phi_{31}) - \frac{A_1^2 G_{33}}{2} \cos[2wt + \phi_2 + \phi_1]$$

$$\begin{aligned}\dot{I}_6 &= C_1 \dot{I}_5 \\ &= \frac{A_1^2 G_{33}}{2} \cos(\phi_{31}) - \frac{C_1 A_1^2 G_{33}}{2} [\cos(2\omega t + \phi_1)]\end{aligned}$$

$$\begin{aligned}\dot{I}_7 &= \dot{I}_3 + \dot{I}_6 \\ &= \frac{A_1^2 C_2}{2} - \frac{A_1^2 C_2}{2} \cos(2\omega t + 2\phi_1) + \frac{A_1^2 G_{33} C_1}{2} \cos(\phi_{31}) \\ &\quad - \frac{A_1^2 G_{33} C_1}{2} \cos(2\omega t + \phi_1 + \phi_{31}) \\ &= \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos(\phi_{31}) \\ &= \frac{A_1^2 C_2}{2} \cos(2\omega t + 2\phi_1) - \frac{A_1^2 G_{33} C_1}{2} \cos(2\omega t + \phi_1 + \phi_{31})\end{aligned}$$

USE MY FORMULA

$$\begin{aligned}B_1 \cos(2\omega t + \psi_1) + B_2 \cos(2\omega t + \psi_2) \\ = \gamma \cos(2\omega t + \psi_1 - \theta)\end{aligned}$$

USE WHICH

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta = \tan^{-1} \frac{q}{P}$$

$$P = B_1 + B_2 \cos(\psi_2 - \psi_1)$$

$$q = B_2 \sin(\psi_1 - \psi_2)$$

$$\sin(\psi) = \cos\left(\psi - \frac{\pi}{2}\right)$$

$$\sin\left(\psi + \frac{\pi}{2}\right) = \cos(\psi)$$

$$\frac{A_1^2 C_2}{2} \cos(2\omega t + 2\phi_1) + \frac{A_1^2 G_{33} C_1}{2} \cos(2\omega t + \phi_1 + \phi_2)$$

$$= \gamma \cos(2\omega t + 2\phi_1 - \theta)$$

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta = \tan^{-1} \frac{P}{q}$$

$$P = \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos(\phi_1 + \phi_2 - 2\phi_1) =$$

$$= \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos[\phi_2 + \phi_1]$$

$$\Rightarrow P = \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos[\phi_{31}]$$

$$q = \frac{A_1^2 G_{33} C_1}{2} \sin[2\phi_1 - \phi_1 - \phi_2]$$

$$= \frac{A_1^2 G_{33} C_1}{2} \sin[\phi_1 - \phi_2]$$

$$= \frac{A_1^2 G_{33} C_1}{2} \sin[-\phi_{31}]$$

$$q = -\frac{A_1^2 G_{33} C_1}{2} \sin[\phi_{31}]$$

$$I_7 = \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos(\phi_{31}) - \cos(2\omega t + 2\phi_1 - \theta)$$

$$\gamma = \sqrt{p^2 + q^2}$$

$$\theta = \tan^{-1} \frac{q}{P}$$

$$P = \frac{A_1^2 C_2}{2} + \frac{A_1^2 G_{33} C_1}{2} \cos(\phi_{31})$$

$$= \frac{A_1^2}{2} [C_2 + G_{33} C_1 \cos(\phi_{31})]$$

$$q = \frac{A_1^2}{2} G_{33} C_1 \sin(\phi_{31})$$

$$I_7 = A_1 + B_7 \cos(2\omega t + \phi_1)$$

USE WHICH:

$$A_7 = \frac{A_1^2}{2} [C_2 + G_{33} C_1 \cos(\phi_{31})]$$

$$B_7 = -\gamma_7$$

$$\phi_7 = 2\phi_1 - \phi_1 \rightarrow \phi_7$$

$$\begin{aligned} \dot{I}_8 &= A_7 F_6(S=0) + B_7 |F_6(S=J2w)| \cos(2wt + \phi_1 + ?? F_6(S=J2w)) \\ &= A_8 + B_8 \cos(2wt + \phi_8) \end{aligned}$$

$$A_8 = A_7 F_6(S=0)$$

$$B_8 = B_7 |F_6(S=J2w)|$$

$$\phi_8 = \phi_7 + ?? F_6(S=J2w)$$

$$\begin{aligned} \dot{I}_{10} &= A_8 F_3(S=0) + B_8 |F_3(S=J2w) \cos(2wt + \phi_8 + ?? F_3(S=J2w))| \\ &= A_{10} + B_{10} \cos(2wt + \phi_{10}) \end{aligned}$$

$$A_{10} = A_8 F_3(S=0)$$

$$B_{10} = B_8 |F_3(S=J2w)|$$

$$\phi_{10} = \phi_8 + ?? F_3(S=J2w)$$

$$C_{11} = C_4 C_7$$

$$= C_4 A_7 + C_4 B_7 \cos(2wt + \phi_1)$$

$$\dot{I}_{12} = \dot{I}_{10} + \dot{I}_{11}$$

$$= A_{10} + B_{10} \cos(2wt + \phi_{10}) + C_4 A_7 + C_4 B_7 \cos(2wt + \phi_1)$$

$$= A_{10} + C_4 A_7 + B_{10} \cos(2wt + \phi_{10}) + C_4 B_7 \cos(2wt + \phi_1)$$

$$B_{10} \cos(2wt + \phi_{10}) + C_4 B_7 \cos(2wt + \phi_1)$$

$$= \gamma \cos[2wt + \phi_{10} - \phi]$$

MOW USE CALCULATE THE H_{3X} TERMS

$$H_{33} = C_7 \dot{I}_1 \dot{I}_{12}$$

$$\begin{aligned} H_{33} &= C_7 A_1 \sin(\omega t + \phi_1) [A_{12} + B_{12} \cos(2\omega t + \phi_{12})] \\ &= C_7 A_1 A_{12} \sin(\omega t + \phi_1) + C_7 A_1 B_{12} \sin(\omega t + \phi_1) \cos(2\omega t + \phi_{12}) \end{aligned}$$

$$\begin{aligned} &\sin(\omega t + \phi_1) \cos(2\omega t + \phi_{12}) \\ &= \frac{1}{2} \sin(3\omega t + \phi_1 + \phi_{12}) + \frac{1}{2} \sin(-\omega t + \phi_1 - \phi_{12}) \\ &= \frac{1}{2} \sin(3\omega t + \phi_1 + \phi_{12}) - \frac{1}{2} \sin(\omega t + \phi_{12} - \phi_1) \end{aligned}$$

$$\begin{aligned} H_{33} &= C_7 A_1 A_{12} \sin(\omega t + \phi_1) + \frac{1}{2} C_7 A_1 B_{12} \sin(3\omega t + \phi_1 + \phi_{12}) \\ &= -\frac{1}{2} C_7 A_1 B_{12} \sin(\omega t + \phi_{12} - \phi_1) \\ &= C_7 A_1 A_{12} \sin(\omega t + \phi_1) - \frac{1}{2} C_7 A_1 B_{12} \sin(\omega t + \phi_{12} - \phi_1) \\ &\quad + \frac{1}{2} C_7 A_1 B_{12} \sin(3\omega t + \phi_1 + \phi_{12}) \end{aligned}$$

$$\sin[\psi] = \cos\left[\psi - \frac{\pi}{2}\right]$$

$$\begin{aligned} H_{33} &= C_7 A_1 A_{12} \cos\left[\omega t + \phi_1 - \frac{\pi}{2}\right] - \frac{1}{2} C_7 A_1 B_{12} \cos\left[\omega t + \phi_{12} - \phi_1 - \frac{\pi}{2}\right] \\ &\quad + \frac{1}{2} C_7 A_1 B_{12} \cos\left[3\omega t + \phi_1 + \phi_{12} - \frac{\pi}{2}\right] \end{aligned}$$

NO MOW MY FORMULA OR PAGE 2

$$\begin{aligned} & C_7 A_1 A_{12} \cos\left[wt + \phi_1 - \frac{\pi}{2}\right] - \frac{1}{2} C_7 A_1 B_{12} \cos\left[wt + \phi_{12} - \phi_1 - \frac{\pi}{2}\right] \\ &= C_7 A_1 A_{12} \cos\left[wt + \phi_1 - \frac{\pi}{2}\right] + \frac{1}{2} C_7 A_1 B_{12} \cos\left[wt + \phi_{12} - \phi_1 + \frac{\pi}{2}\right] \\ &= \gamma \cos\left[wt + \phi_1 - \frac{\pi}{2} - \theta\right] \end{aligned}$$

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta = \tan^{-1} \frac{q}{P}$$

$$\begin{aligned} P &= C_7 A_1 A_{12} + \frac{1}{2} C_7 A_1 B_{12} \cos\left[\phi_{12} - \phi_1 + \frac{\pi}{2} - \phi_1 + \frac{\pi}{2}\right] \\ &= C_7 A_1 A_{12} + \frac{1}{2} C_7 A_1 B_{12} \cos[\phi_{12} - 2\phi_1 + \pi] \\ &= C_7 A_1 A_{12} - \frac{1}{2} C_7 A_1 B_{12} \cos[\phi_{12} - 2\phi_1] \end{aligned}$$

$$\begin{aligned} q &= \frac{1}{2} C_7 A_1 B_{12} \sin[2\phi_1 - \phi_{12} - \pi] \\ &= -\frac{1}{2} C_7 A_1 B_{12} \sin[2\phi_1 - \phi_{12}] \\ &= \frac{1}{2} C_7 A_1 B_{12} \sin[\phi_{12} - 2\phi_1] \end{aligned}$$

$$H_{33} = \gamma \cos\left[wt + \phi_1 - \frac{\pi}{2} + \theta\right] + \frac{1}{2} C_7 A_1 B_{12} \cos\left[3wt + \phi_1 + \phi_{12} - \frac{\pi}{2}\right]$$

$$H_{33} = H_{331} + H_{333}$$

$$\begin{aligned} H_{331} &= \gamma \cos\left[wt + \phi_1 - \frac{\pi}{2} - \theta\right] \\ &= m_{331} \cos[wt + \phi_{331}] \end{aligned}$$

$$m_{331} = \gamma$$

$$\phi_{331} = \phi_1 - \frac{\pi}{2} - \theta$$

$$H_{333} = C_7 A_1 B_{12} \cos\left[3wt + \phi_1 + \phi_{12} - \frac{\pi}{2}\right]$$

$$= m_{333} \cos[3\omega t + \phi_{333}]$$

$$\phi_{333} = \phi_1 + \phi_{12} - \frac{\pi}{2}$$

$$m_{333} = \frac{1}{2} C_7 A_1 B_{12}$$

$$H_{31} = C_5 \dot{I}_1 \dot{I}_8$$

$$= C_5 A_1 A_8 \sin(\omega t + \phi_1) [A_8 + B_8 \cos(2\omega t + \phi_8)]$$

$$= C_5 A_1 A_8 \sin[(\omega t + \phi_1)]$$

$$+ C_5 A_1 B_8 \sin(\omega t + \phi_1) \cos(2\omega t + \phi_8)$$

$$\sin(\omega t + \phi_1) \cos(2\omega t + \phi_8) = \frac{1}{2} \sin(3\omega t + \phi_1 + \phi_8) - \frac{1}{2} \sin(\omega t + \phi_8 - \phi_1)$$

$$H_{31} = C_5 A_1 A_8 \sin(\omega t + \phi_1) - \frac{1}{2} C_5 A_1 B_8 \sin(\omega t + \phi_8 - \phi_1)$$

$$+ \frac{1}{2} C_5 A_1 B_8 \sin(3\omega t + \phi_1 + \phi_8)$$

$$\sin(\psi) = \cos\left(\psi - \frac{\pi}{2}\right)$$

$$H_{31} = C_5 A_1 A_8 \cos(\omega t + \phi_1 - \frac{\pi}{2}) - \frac{1}{2} C_5 A_1 B_8 \cos(\omega t + \phi_8 - \phi_1 - \frac{\pi}{2})$$

$$+ \frac{1}{2} C_5 A_1 B_8 \cos(3\omega t + \phi_1 + \phi_8 - \frac{\pi}{2})$$

MOW USING MY FORMULA ON PAGE 2

$$C_5 A_1 A_8 \cos\left[\omega t + \phi_1 - \frac{\pi}{2}\right] - \frac{1}{2} C_5 A_1 B_8 \cos\left[\omega t + \phi_8 - \phi_1 - \frac{\pi}{2}\right]$$

$$= C_5 A_1 A_8 \cos\left[\omega t + \phi_1 - \frac{\pi}{2}\right] + \frac{1}{2} C_5 A_1 B_8 \cos\left[\omega t + \phi_8 - \phi_1 + \frac{\pi}{2}\right]$$

$$= \gamma \cos\left[\omega t + \phi_1 - \frac{\pi}{2} - \theta\right]$$

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta_{m6} = \tan^{-1} \frac{q}{P}$$

$$P = C_5 A_1 A_8 + \frac{1}{2} C_5 A_1 B_8 \cos\left(\phi_8 - \phi_1 + \frac{\pi}{2} - \phi_1 + \frac{\pi}{2}\right)$$

$$= C_5 A_1 A_8 + \frac{1}{2} C_5 A_1 B_8 \cos(\phi_8 - 2\phi_1 + \pi)$$

$$= C_5 A_1 A_8 - \frac{1}{2} C_5 A_1 B_8 \cos(\phi_8 - 2\phi_1)$$

$$q = \frac{1}{2} C_5 A_1 A_8 \sin(2\phi_1 - \phi_8 - \pi)$$

$$= \frac{1}{2} C_5 A_1 B_8 \sin(\phi_8 - 2\phi_1)$$

$$H_{31} = \gamma \cos\left[wt + \phi_1 - \frac{\pi}{2} - \theta\right] + \frac{1}{2} C_5 A_1 B_8 \cos\left(3wt + \phi_1 + \phi_8 - \frac{\pi}{2}\right)$$

$$= H_{311} + H_{313}$$

$$H_{311} = m_{61} \cos[wt + \phi_{61}]$$

$$m_{61} = \gamma$$

$$\phi_{61} = \phi_1 - \frac{\pi}{2} - \theta$$

$$H_{313} = m_{63} \cos[3wt + \phi_{63}]$$

$$m_{63} = \frac{1}{2} C_5 A_1 B_8$$

$$\phi_{63} = \phi_1 + \phi_8 - \frac{\pi}{2}$$

$$H_{32} = C_6 \dot{I}_4 \dot{I}_8$$

$$= C_6 A_1 G_{33} \sin(wt + \phi_2) [A_8 + B_8 \cos(2wt + \phi_8)]$$

$$C_5 \rightarrow C_6 G_{33}$$

$$H_{32} = \gamma \cos\left[wt + \phi_1 - \frac{\pi}{2} - \theta\right] + \frac{1}{2} C_6 G_{33} A_1 B_8 \cos\left[3wt + \phi_1 + \phi_8 - \frac{\pi}{2}\right]$$

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta = \tan^{-1} \frac{q}{P}$$

$$P = C_6 G_{33} A_1 A_8 - \frac{1}{2} C_6 G_{33} A_1 B_8 \cos(\phi_8 - 2\phi_1)$$

$$q = \frac{1}{2} C_6 G_{33} A_1 B_8 - \sin(\phi_8 - 2\phi_1)$$

$$H_{32} = H_{321} + H_{323}$$

$$H_{321} = m_{71} \cos[wt + \phi_{71}]$$

$$m_{71} = \gamma$$

$$\phi_{71} = \phi_1 - \frac{\pi}{2} - \theta$$

$$H_{323} = m_{73} \cos[3wt + \phi_{73}]$$

$$m_{73} = \frac{1}{2} C_6 G_{33} A_1 B_8$$

$$\phi_{73} = \phi_1 + \phi_8 - \frac{\pi}{2}$$

$$H_{34} = C_8 \dot{I}_2 \dot{I}_4$$

$$\dot{I}_2 = \frac{A_1^2}{2} - \frac{A_1^2}{2} \cos(2wt + 2\phi_1)$$

$$\dot{I}_4 = A_1 G_{33} \sin(wt + \phi_2)$$

$$H_{34} = \frac{A_1^2}{2} [1 - \cos(2wt + 2\phi_1)] A_1 G_{33} \sin(wt + \phi_2)$$

$$= \frac{1}{2} A_1^3 G_{33} [1 - \cos(2wt + 2\phi_1)] [\sin(wt + \phi_2)]$$

$$C_5 A_1 \rightarrow \frac{1}{2} A_1^3 G_{33}$$

$$\phi_1 \rightarrow \phi_2$$

$$\phi_8 \rightarrow 2\phi_1$$

$$A_8 = 1$$

$$B_8 = -1$$

$$H_{34} = \gamma \cos\left(wt + \phi_2 - \frac{\pi}{2} - \theta\right) + \frac{1}{2} \frac{1}{2} A_1^3 G_{33} (-1) \cos\left[3wt + \phi_2 + 2\phi_1 - \frac{\pi}{2}\right]$$

$$= \gamma \cos\left(wt + \phi_2 - \frac{\pi}{2} - \theta\right) - \frac{1}{4} A_1^3 G_{33} \cos\left[3wt + \phi_2 + 2\phi_1 - \frac{\pi}{2}\right]$$

$$\gamma \cos\left(wt + \phi_2 - \frac{\pi}{2} - \theta \right) + \frac{1}{4} A_1^3 G_{33} \cos\left[3wt + \phi_1 + 2\phi_1 + \frac{\pi}{2} \right]$$

$$\gamma = \sqrt{P^2 + q^2}$$

$$\theta = \tan^{-1} \frac{q}{P}$$

$$P = \frac{1}{2} A_1^3 G_{33} + \frac{1}{2} \frac{1}{2} A_1^3 G_{33} \cos[2\phi_1 - 2\phi_2]$$

$$= \frac{1}{2} A_1^3 G_{33} + \frac{1}{4} A_1^3 G_{33} \cos[2(\phi_1 - \phi_2)]$$

$$q = \frac{1}{2} \frac{1}{2} A_1^3 G_{33} (-1) \sin[2\phi_1 - 2\phi_2]$$

$$= -\frac{1}{4} A_1^3 G_{33} \sin[2(\phi_1 - \phi_2)]$$

$$= +\frac{1}{4} A_1^3 G_{33} \sin[2(\phi_2 - \phi_1)]$$

$$H_{34} = H_{341} + H_{343}$$

$$H_{341} = m_{81} \cos[wt + \phi_{81}]$$

$$m_{81} = \gamma$$

$$\phi_{81} = \phi_2 - \frac{\pi}{2} - \theta$$

$$m_{83} = \frac{1}{4} A_1^3 G_{33}$$

$$\phi_{83} = \phi_2 + 2\phi_1 + \frac{\pi}{2}$$

WE ARE MOW AT

$$H_{33} = H_{331} + H_{333}$$

$$H_{31} = H_{311} + H_{313}$$

$$H_{32} = H_{321} + H_{323}$$

$$H_{34} = H_{341} + H_{343}$$

$$(_{331} + H_{311}) = m_{331} \cos(wt + \phi_{331}) + m_{61} \cos(wt + \phi_{m61})$$

$$\gamma_{a1} = \cos(wt + \phi_{331} - \theta_{a1})$$

$$\gamma_{200} \rightarrow \gamma_{a1} = \sqrt{P_{a1}^2 + q_{a1}^2}$$

$$\theta_{a1} = \tan^{-1} \frac{q_a}{p_a} \leftarrow \theta_{200}$$

$$P_{200} \rightarrow P_{a1} = m_{331} + m_{61} \cos(\theta_{61} - \theta_{331})$$

$$q_{200} \rightarrow q_{a1} = m_{61} \sin(\theta_{331} - \theta_{61})$$

$$\begin{aligned} H_{331} + H_{311} + H_{321} &= \gamma_{a1} \cos(wt + \phi_{331} - \theta_{a1}) + m_{71} \cos(wt + \phi_{71}) \\ &= \gamma_{b1} \cos(wt + \phi_{331} - \theta_{a1} - \theta_{b1}) \end{aligned}$$

$$\gamma_{201} \rightarrow \gamma_{b1} = \sqrt{P_{b1}^2 + q_{a1}^2}$$

$$\theta_{b1} = \tan^{-1} \frac{q_{b1}}{p_{b1}} \leftarrow \theta_{201}$$

$$P_{201} \rightarrow P_{b1} = \gamma_{a1} + m_{61} \cos(\theta_{71} - \theta_{331} + \theta_{a1})$$

$$q_{201} \rightarrow q_{b1} = m_{71} \sin(\theta_{331} - \theta_{a1} - \theta_{71})$$

$$S_{ma1} = (H_{331} + H_{321} + H_{321}) + H_{341}$$

$$S_{ma1} = \gamma_{b1} \cos(wt + \theta_{331} - \theta_{a1} - \theta_{b1}) + m_{81} \cos(wt + \theta_{81})$$

$$S_{ma1} = \gamma_{b1} \cos(wt + \theta_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1})$$

$$\gamma_{c1} = \sqrt{P_{c1}^2 + q_{c1}^2}$$

$$\theta_{c1} = \tan^{-1} \frac{q_{c1}}{p_{c1}}$$

$$P_{c1} = \gamma_{b1} + m_{81} \cos(\phi_{81} - \phi_{331} + \theta_{a1} + \theta_{b1})$$

$$q_{c1} = m_{81} \sin(\phi_{331} - \theta_{a1} - \theta_{b1} - \phi_{81})$$

$$S_{m23} = H_{333} + H_{313} + H_{323} + H_{343}$$

17.SAYFA

$$H_{333} + H_{313} = m_{333} \cos(3wt + \phi_{333}) + m_{63} \cos(3wt + \phi_{63})$$

$$= \gamma_{a3} \cos(3wt + \phi_{333} - \phi_{a3})$$

$$\gamma_{230} \rightarrow \gamma_{a3} = \sqrt{P_{a3}^2 + q_{a3}^2}$$

$$\theta_{a3} = \tan^{-1} \frac{q_{a3}}{p_{a3}}$$

$$P_{230} \rightarrow P_{a3} = m_{333} + m_{63} \cos(\phi_{63} - \phi_{333})$$

$$q_{230} \rightarrow q_{a3} = m_{63} \sin(\phi_{333} - \phi_{63})$$

$$(H_{333} + H_{313}) + H_{323} = \gamma_{a3} \cos(3wt + \phi_{333} - \phi_{a3}) + m_{63} \cos(3wt + \phi_{73})$$

$$= \gamma_{b3} \cos(3wt + \phi_{333} - \phi_{a3} - \theta_{b3})$$

$$\gamma_{231} \rightarrow \gamma_{b3} = \sqrt{P_{b3}^2 + q_{b3}^2}$$

$$\theta_{b3} = \tan^{-1} \frac{q_{b3}^2}{P_{b3}^2} \leftarrow \theta_{231}$$

$$P_{231} \rightarrow P_{b3} = \gamma_{a3} + m_{73} \cos(\phi_{73} - \phi_{333} + \theta_{a3})$$

$$q_{231} \rightarrow q_{b3} = m_{73} \sin(\phi_{333} - \phi_{a3} - \phi_{73})$$

$$S_{ma3} = (H_{333} + H_{313} + H_{323}) + H_{343}$$

$$= \gamma_{b3} \cos(3wt + \phi_{333} - \theta_{a3} - \theta_{b3}) + m_{83} \cos(3wt + \phi_{83})$$

$$= \gamma_{c3} \cos(3wt + \phi_{333} - \theta_{a3} - \theta_{b3} - \theta_{c3})$$

$$\gamma_{232} = \gamma_{c3} = \sqrt{P_{c3}^2 + q_{c3}^2}$$

$$\theta_{c3} = \tan^{-1} \frac{q_{c3}^2}{P_{c3}^2} \leftarrow \theta_{232}$$

$$P_{232} \rightarrow P_{c3} = \gamma_{b3} + m_{83} \cos(\phi_{83} - \phi_{333} + \theta_{a3} + \theta_{b3})$$

$$q_{232} \rightarrow q_{c3} = m_{83} \sin(\phi_{333} - \theta_{a3} - \theta_{b3} - \phi_{83})$$

$$S_{m20} = S_{m21} + S_{m23}$$

$$S_{m21} = \gamma_{c1} \cos(wt + \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1})$$

$$S_{m23} = \gamma_{c3} \cos(3wt + \phi_{331} - \theta_{a3} - \theta_{b3} - \theta_{c3})$$

$$F10 = F6.R1$$

$$PH_{31} = PH_{311} + PH_{313}$$

$$PH_{311} = |F10(S = Jw)| \gamma_{c1} \cos(wt + \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1}) + \kappa F10(S = Jw)$$

$$PH_{313} = |F10(S = J3w)| \gamma_{c3} \cos(3wt + \phi_{331} - \theta_{a3} - \theta_{b3} - \theta_{c3}) + \kappa F10(S = J3w)$$

$$\dot{I}_8 = A_8 + B_8 \cos(2wt + \phi_8)$$

$$\dot{I}_{14} = -A_8 K q_0 + B_8 K q_0 \cos[2wt + \phi_8 + \pi - 2w]$$

$$\begin{aligned}
I_{15} &= \dot{I}_{14} + \dot{I} \\
&= \dot{I}_{14} + \mu \sin(wt) \\
&= -A_8 K q_0 + \mu \sin(wt) + B_8 K q_0 \cos[2wt + \phi_8 + \pi - 2w] \\
I_{16} &= -A_8 K q_0 Z_1(S=0) + \mu |Z_1(S=Jw)| \sin[wt + \kappa Z_1(S=jw)] \\
&+ A_8 K q_0 \cos[2wt + \phi_8 + \pi - 2w + \kappa Z_1(S=j2w)] \\
&= A_{16} + C_{16} \sin[wt + \Phi_{11}] + B_{16} \cos[2wt + \phi_{162}] \\
A_{16} &= -A_8 K q_0 Z_1(S=0) \\
C_{16} &= \mu |Z_1(S=Jw)| \\
\phi_{161} &\rightarrow \phi_{11} = \kappa Z_1(S=Jw) \rightarrow \phi_{11} \\
B_{16} &= B_8 K q_0 |Z_1(S=j2w)| \\
\phi_{162} &= \phi_8 + \pi - 2w + \kappa Z_1(S=J2w) \\
\dot{I}_{12} &= \dot{I}_{16}^2 = [A_{16} + C_{16} \sin(wt + \phi_{11}) + B_{16} \cos(2wt + \phi_{162})] \\
&= 2A_{16}C_{16} \sin(wt + \phi_{11}) + 2C_{16}B_{16} \sin(wt + \phi_{11}) \cos(2wt + \phi_{162}) \\
&= 2A_{16}C_{16} \sin(wt + \phi_{11}) + C_{16}B_{16} \sin(3wt + \phi_{11} + \phi_{162}) \\
&+ C_{16}B_{16} \sin(-wt + \phi_{11} - \phi_{162}) \\
&= 2A_{16}C_{16} \sin(wt + \phi_{11}) - C_{16}B_{16} \sin(wt - \phi_{11} + \phi_{162}) + C_{16}B_{16} \sin(3wt + \phi_{11} + \phi_{162})
\end{aligned}$$

$$2A_{16}C_{16} \cos(wt + \phi_{11} - \frac{\pi}{2}) + C_{16}B_{16} \cos(wt - \phi_{11} + \phi_{162} + \frac{\pi}{2})$$

$$= \gamma_{17} \cos(wt + \phi_{11} - \frac{\pi}{2} - \phi_{17})$$

$$\gamma_{17} = \sqrt{P_{17}^2 + q_{17}^2}$$

$$\theta_{17} = \tan^{-1} \frac{q_{17}}{P_{17}}$$

$$P_{17} = 2A_{16}C_{16} + C_{16}B_{16} \cos(\phi_{162} - \phi_{11} + \frac{\pi}{2} - \phi_{11} + \frac{\pi}{2})$$

$$= 2A_{16}C_{16} - C_{16}B_{16} \cos(\phi_{162} - 2\phi_{11})$$

$$q_{17} = C_{16}B_{16} \sin(\phi_{11} - \frac{\pi}{2} - \phi_{162} + \phi_{11} - \frac{\pi}{2})$$

$$= -C_{16}B_{16} \sin(2\phi_{11} - \phi_{162})$$

$$\begin{aligned} \dot{I}_{17} &= \gamma_{17} \cos(wt + \phi_{11} - \frac{\pi}{2} - \theta_{17}) + C_{16} B_{16} \sin(3wt + \phi_{11} + \phi_{162}) \\ &= \gamma_{17} \cos(wt + \psi_{171}) + C_{16} B_{16} \sin(3wt + \psi_{173}) \end{aligned}$$

$$\psi_{171} = \phi_{11} - \frac{\pi}{2} - \theta_{17}$$

$$\psi_{173} = \phi_{11} + \phi_{162}$$

$$\begin{aligned} \dot{I}_{18} &= C_2 \dot{I}_{17} \\ &= C_2 \gamma_{17} \cos(wt + \psi_{171}) + C_2 C_{16} B_{16} \sin(3[wt + \psi_{173}]) \end{aligned}$$

$$\dot{I}_{180} = A_{18} \cos(wt + \psi_{181}) + B_{18} \sin(3wt + \psi_{183})$$

$$A_{18} = C_2 \gamma_{17}$$

$$B_{18} = C_2 C_{16} B_{16}$$

$$\begin{aligned} \dot{I}_{19} &= A_{16} F_3(S=0) + C_{16} |F_3(S=jw)| \sin[wt + \phi_{11} + \kappa F_3(S=jw)] \\ &+ B_{16} |F_3(S=j2w)| \cos[2wt + \phi_{162} + \kappa F_3(S=j2w)] \end{aligned}$$

$$\dot{I}_{19} = A_{19} + C_{19} \sin[wt + \psi_{191}] + B_{19} \cos[2wt + \psi_{192}]$$

$$A_{19} = A_{16} F_3(S=0) \rightarrow F_{30}$$

$$C_{19} = C_{16} |F_3(S=jw)| \rightarrow F_{31}$$

$$B_{19} = B_{16} |F_3(S=j2w)| \rightarrow F_{32}$$

$$\psi_{191} = \phi_{11} + \kappa F_3(S=jw) \rightarrow F_{31} \rightarrow \phi_{31}$$

$$\psi_{192} = \phi_{162} + \kappa F_3(S=j2w) F_{32} \rightarrow \phi_{32}$$

$$\dot{I}_{20} = \dot{I}_{16} \cdot \dot{I}_{19}$$

$$\dot{I}_{20} = [A_{19} + C_{19} \sin[wt + \psi_{191}] + B_{19} \cos[2wt + \psi_{192}]]$$

$$[A_{16} + C_{16} \sin[wt + \phi_{11}] + B_{16} \cos[2wt + \phi_{162}]]$$

AGAIN WE ARE ONLY INTERESTED IN THE HARMONIC TERMS. THEY ARE

$$= A_{19} C_{16} \sin(wt + \phi_{11}) + A_{16} C_{19} \sin(wt + \psi_{191}) + C_{19} B_{16} \sin(wt + \psi_{191}).$$

$$\cos(2wt + \phi_{162}) + C_{16} B_{19} \sin(wt + \phi_{11}) \cos(2wt + \psi_{192})$$

$$= A_{19} C_{16} \sin(wt + \phi_{11}) \mp A_{16} C_{19} \sin(wt + \psi_{191}) + \frac{1}{2} C_{19} B_{16} \sin(3wt + \psi_{191} + \phi_{162})$$

$$\begin{aligned}
& + \frac{1}{2} C_{19} B_{16} \sin(-wt + \psi_{191} - \phi_{162}) + \frac{1}{2} C_{16} B_{19} \sin(3wt + \phi_{11} + \psi_{192}) \\
& + \frac{1}{2} C_{16} B_{19} \sin(-wt + \phi_{11} - \psi_{192})
\end{aligned}$$

$$\begin{aligned}
\dot{I}_{20} &= A_{19} C_{16} \sin(wt + \phi_{11}) + A_{16} C_{19} \sin(wt + \psi_{191}) \\
&- \frac{1}{2} C_{19} B_{16} \sin(3wt + \phi_{162} - \psi_{191}) - \frac{1}{2} C_{16} B_{19} \sin(wt - \phi_{11} + \psi_{192}) \\
&+ \frac{1}{2} C_{19} B_{16} \sin(3wt + \phi_{162} + \psi_{191}) + \frac{1}{2} C_{16} B_{19} \sin(3wt + \phi_{11} + \psi_{192}) \\
&A_{19} C_{16} \cos(wt + \phi_{11} - \frac{\pi}{2}) + A_{16} C_{19} \cos(wt + \psi_{191} - \frac{\pi}{2}) \\
&= \gamma_{20a} \cos(wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a})
\end{aligned}$$

$$\gamma_{300} = \gamma_{20a} = \sqrt{P_{20a}^2 + q_{20a}^2}$$

$$\theta_{20a} = \tan^{-1} \frac{q_{20a}}{P_{20a}}$$

$$P_{300} \rightarrow P_{20a} = A_{19} C_{16} + A_{16} C_{19} \cos(\psi_{191} - \frac{\pi}{2} - \phi_{11} + \frac{\pi}{2})$$

$$P_{300} \rightarrow P_{20a} = A_{19} C_{16} + A_{16} C_{19} \cos(\psi_{191} - \phi_{11})$$

$$q_{300} = q_{20a} = A_{16} C_{19} \sin\left[\phi_{11} - \frac{\pi}{2} - \psi_{191} + \frac{\pi}{2}\right]$$

$$q_{20a} = A_{16} C_{19} \sin[\phi_{11} - \psi_{191}]$$

$$= \gamma_{20a} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a}\right] + \frac{1}{2} C_{19} B_{16} \cos\left[wt - \psi_{191} + \phi_{162} + \frac{\pi}{2}\right]$$

$$\gamma_{30b} = \gamma_{20b} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b}\right]$$

$$\gamma_{20b} = \sqrt{P_{20b}^2 + q_{20b}^2}$$

$$\theta_{20b} = \tan^{-1} \frac{q_{20b}}{P_{20b}}$$

$$P_{30b} = P_{20b} = \gamma_{20a} + \frac{1}{2} C_{19} B_{16} \cos\left[-\psi_{191} + \phi_{162} + \frac{\pi}{2} - \phi_{11} + \frac{\pi}{2} + \theta_{20a}\right]$$

$$\begin{aligned}
&= \gamma_{20a} - \frac{1}{2} C_{19} B_{16} \cos[\phi_{162} - \psi_{191} - \phi_{11} + \theta_{20a}] \\
q_{30b} &= q_{20b} = \frac{1}{2} C_{19} B_{16} \sin\left[\phi_{11} - \frac{\pi}{2} - \theta_{20a} + \psi_{191} - \phi_{162} - \frac{\pi}{2}\right] \\
&= +\frac{1}{2} C_{19} B_{16} \sin[\phi_{162} - \psi_{191} - \phi_{11} + \theta_{20a}] \\
\gamma_{301} &= \gamma_{20b} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b}\right] - \frac{1}{2} C_{16} B_{19} \sin[wt - \phi_{11} - \psi_{192}] \\
&= \gamma_{20b} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b}\right] + \frac{1}{2} C_{16} B_{19} \sin\left[wt - \phi_{11} + \phi_{192} + \frac{\pi}{2}\right] \\
\gamma_{302} &= \gamma_{201} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b} - \theta_{20c}\right] \\
\gamma_{302} &\rightarrow \gamma_{201} = \sqrt{P_{20c}^2 + q_{20c}^2} \\
\theta_{20c} &= \tan^{-1} \frac{q_{20c}}{P_{20c}} \\
P_{20c} &= \gamma_{20b} + \frac{1}{2} C_{16} B_{19} \cos\left[-\phi_{11} + \psi_{192} + \frac{\pi}{2} - \phi_{11} + \frac{\pi}{2} + \theta_{20a} + \theta_{20b}\right] \\
&= \gamma_{20b} - \frac{1}{2} C_{16} B_{19} \cos[+\psi_{192} - 2\phi_{11} + \theta_{20a} + \theta_{20b}] \\
q_{20c} &= \frac{1}{2} C_{16} B_{19} \sin\left[\phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b} + \phi_{11} - \psi_{192} - \frac{\pi}{2}\right] \\
&= -\frac{1}{2} C_{16} B_{19} \sin[2\phi_{11} - \psi_{192} - \theta_{20a} - \theta_{20b}] \\
&= +\frac{1}{2} C_{16} B_{19} \sin[\psi_{192} - 2\phi_{11} + \theta_{20a} + \theta_{20b}] \\
&= \frac{1}{2} C_{19} B_{16} \sin[3wt + \psi_{192} + \phi_{162}] + \frac{1}{2} C_{16} B_{19} \sin[3wt + \psi_{192} + \phi_{11}] \\
&= \frac{1}{2} \left[C_{19} B_{16} \cos\left[3wt + \psi_{192} + \phi_{162} - \frac{\pi}{2}\right] + C_{16} B_{19} \cos\left[3wt + \psi_{192} + \phi_{11} - \frac{\pi}{2}\right] \right] \\
&= \gamma_{203} \cos\left[3wt + \psi_{192} + \phi_{162} - \frac{\pi}{2} - \theta_{203}\right] \\
\gamma_{203} &= \frac{1}{2} \sqrt{P_{203}^2 + q_{203}^2}
\end{aligned}$$

$$\begin{aligned}
P_{203} &= C_{19}B_{16} + C_{16}B_{19} \cos\left[\phi_{11} + \psi_{192} - \frac{\pi}{2} - \psi_{191} - \phi_{162} + \frac{\pi}{2}\right] \\
&= C_{19}B_{16} + C_{16}B_{19} \cos[\phi_{11} + \psi_{192} - \psi_{191} - \phi_{162}] \\
q_{203} &= C_{16}B_{19} \sin\left[\psi_{191} + \phi_{162} - \frac{\pi}{2} - \phi_{11} - \psi_{192} + \frac{\pi}{2}\right] \\
&= C_{16}B_{19} \sin[\psi_{191} + \phi_{162} - \phi_{11} - \psi_{192}]
\end{aligned}$$

$$\begin{aligned}
\dot{I}_{20} &= \gamma_{201} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b} - \theta_{20c}\right] \\
&+ \gamma_{201} \cos\left[3wt + \psi_{191} + \phi_{162} - \frac{\pi}{2} - \theta_{203}\right] \\
\dot{I}_{21} &= C_1 \dot{I}_{20} \\
&= C_1 \gamma_{201} \cos\left[wt + \phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b} - \theta_{20c}\right] \\
&+ C_1 \gamma_{203} \cos\left[3wt + \psi_{191} + \phi_{162} - \frac{\pi}{2} - \theta_{203}\right] \\
\dot{I}_{22} &= \dot{I}_{18} + \dot{I}_{21} \\
&= C_2 \gamma_{17} \cos[wt + \psi_{191}] + C_2 C_{16} B_{16} \sin[3wt + \psi_{173}] + C_1 \gamma_{201} \cos[wt + \psi_{211}] \\
&+ C_1 \gamma_{203} \cos[3wt + \psi_{213}] \\
\psi_{211} &= \left[\phi_{11} - \frac{\pi}{2} - \theta_{20a} - \theta_{20b} - \theta_{20c}\right] \\
\psi_{213} &= \left[\psi_{191} + \phi_{162} - \frac{\pi}{2} - \theta_{203}\right]
\end{aligned}$$

SO FIRST HARMONIC??

$$\begin{aligned}
&= C_2 \gamma_{17} \cos[wt + \psi_{171}] + C_2 \gamma_{201} \cos[wt + \psi_{211}] \\
&= C_{221} \cos[wt + \psi_{171} - \theta_{221}] \\
\gamma_{221} &= \sqrt{P_{221}^2 + q_{221}^2} \\
P_{221} &= C_2 \gamma_{17} + C_1 \gamma_{201} \cos[\psi_{221} - \psi_{171}] \\
q_{221} &= C_1 \gamma_{201} \sin[\psi_{171} - \psi_{211}]
\end{aligned}$$

AND THE 3 ND HARMANIC ...??

$$\begin{aligned}
&= C_2 C_{16} B_{16} \sin[3wt + \psi_{173}] + C_1 \gamma_{203} \cos[3wt + \psi_{213}] \\
&= C_2 C_{16} B_{16} \cos\left[3wt + \psi_{173} - \frac{\pi}{2}\right] + C_1 \gamma_{203} \cos[3wt + \psi_{213}] \\
&= \gamma_{223} \cos\left[3wt + \psi_{173} - \frac{\pi}{2} - \theta_{223}\right] \\
\gamma_{223} &= \sqrt{P_{223}^2 + q_{223}^2} \\
\theta_{223} &= \tan^{-1} \frac{q_{223}}{P_{223}} \\
P_{223} &= C_2 C_{16} B_{16} + C_1 \gamma_{203} \cos\left[\psi_{213} - \psi_{173} + \frac{\pi}{2}\right] \\
q_{223} &= C_1 \gamma_{203} \sin\left[\psi_{173} - \frac{\pi}{2} - \psi_{213}\right] \\
\dot{I}_{22} &= \gamma_{221} \cos[wt + \psi_{171} - \theta_{221}] + \gamma_{223} \cos\left[3wt + \psi_{173} - \frac{\pi}{2} - \theta_{223}\right] \\
PH_{32} &= \gamma_{221} |F_{10}(S = jw)| \cos[wt + \psi_{171} - \theta_{221} + \kappa F_{10}(S = jw)] \\
&+ \gamma_{223} |F_{10}(S = j3w)| \cos\left[3wt + \psi_{173} - \frac{\pi}{2} - \theta_{223} + \kappa F_{10}(S = j3w)\right] \\
&= PH_{321} + PH_{321} \\
PH_{321} &= \gamma_{221} |F_{10}(S = jw)| \cos[wt + \psi_{171} - \theta_{221} + \kappa F_{10}(S = jw)] \\
PH_{323} &= \gamma_{223} |F_{10}(S = j3w)| \cos\left[3wt + \psi_{173} - \frac{\pi}{2} - \theta_{223} + \kappa F_{10}(S = j3w)\right] \\
P_3 &= PH_{31} + PH_{32} = PH_{311} + PH_{313} + PH_{321} + PH_{323} \\
&= P_{P31} + P_{P32} \\
P_{P31} &= PH_{311} + PH_{321} = |F_{10}(S = j3w)| \gamma_{C1} \cos \\
&[wt + \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1} + \kappa F_{10}(S = jw)] \\
&+ \gamma_{221} |F_{10}(S = j3w)| \cos[wt + \psi_{171} - \theta_{221} + \kappa F_{10}(S = jw)] \\
&= Z_{31} \cos[wt + \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1} + \kappa F_{10}(S = jw)]
\end{aligned}$$

$$Z_{31} = |F_{10}(S = jw)|\sqrt{P_{31}^2 + q_{31}^2}$$

$$\theta_{p31} = \tan^{-1} \frac{q_{31}}{P_{31}}$$

$$P_{31} = \gamma_{c1} + \gamma_{221} \cos[\psi_{171} - \theta_{221} - \phi_{331} + \theta_{a1} + \theta_{b1} + \theta_{c1}]$$

$$q_{31} = \gamma_{221} \sin[\theta_{221} - \psi_{171} - \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1}]$$

$$P_{p33} = PH_{313} + PH_{323}$$

$$= \gamma_{c3} |F_{10}(S = jw)| \cos[3wt + \phi_{333} - \theta_{a1} - \theta_{b1} - \theta_{c1} + \kappa F_{10}(S = j3w)]$$

$$+ \gamma_{223} |F_{10}(S = j3w)| \cos\left[3wt + \psi_{123} - \frac{\pi}{2} - \theta_{233} + \kappa F_{10}(S = j3w)\right]$$

$$= Z_{33} \cos[3wt + \phi_{333} - \theta_{a3} - \theta_{b3} - \theta_{c3} + \kappa F_{10}(S = j3w) - \theta_{p3}]$$

$$Z_{33} = |F_{10}(S = j3w)|\sqrt{P_{33}^2 + q_{33}^2}$$

$$\theta_{p33} = \tan^{-1} \frac{q_{33}}{P_{33}}$$

$$P_{33} = \gamma_{c3} + \gamma_{223} \cos\left[\psi_{123} - \frac{\pi}{2} - \theta_{233} - \phi_{333} - \theta_{a3} + \theta_{b3} + \theta_{c3}\right]$$

$$= \gamma_{c3} + \gamma_{223} \cos[\psi_{123} - \theta_{233} - \phi_{333} + \theta_{a3} + \theta_{b3} + \theta_{c3}]$$

$$q_{33} = \gamma_{223} \sin\left[\phi_{333} - \theta_{a3} - \theta_{b3} - \theta_{c3} - \psi_{173} + \frac{\pi}{2} + \theta_{223}\right]$$

$$= \gamma_{223} \cos[\phi_{333} - \theta_{a3} - \theta_{b3} - \theta_{c3} - \psi_{173} + \theta_{223}]$$

$$P_3 = P_{p31} + P_{p33}$$

$$P_{p31} = Z_{31} \cos[wt + \phi_{331} - \theta_{a1} - \theta_{b1} - \theta_{c1} + \kappa F_{10}(S = jw) - \theta_{p3}]$$

$$P_{p33} = Z_{33} \cos[3wt + \phi_{333} - \theta_{a3} - \theta_{b3} - \theta_{c3} + \kappa F_{10}(S = j3w) - \theta_{p3}]$$